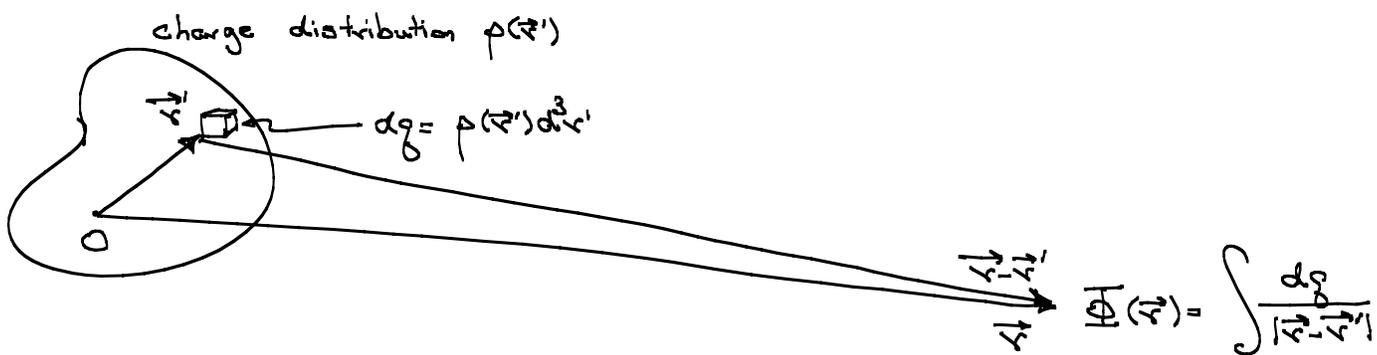


8.4



(a)
$$\frac{1}{|\mathbf{r} - \mathbf{r}'|} = \frac{1}{r} \left[1 - \frac{\mathbf{r} \cdot \mathbf{r}'}{r^2} + \frac{3}{2} \frac{(\mathbf{r} \cdot \mathbf{r}')^2}{r^4} - \dots \right]$$

\leftarrow this term is $O(r'/r)$
 \leftarrow this term is $O((r'/r)^2)$

Expand this using

$$(1+x)^{-1/2} = 1 - \frac{1}{2}x + \frac{3}{8}x^2 - \dots$$

$$= 1 - \frac{1}{2}x + \dots$$

\leftarrow this term is $O(r'/r)$
 \leftarrow this term is $O((r'/r)^2)$

$$\frac{1}{|\mathbf{r} - \mathbf{r}'|} = \frac{1}{r} \left[1 - \frac{\mathbf{r} \cdot \mathbf{r}'}{r^2} + \frac{3}{2} \frac{(\mathbf{r} \cdot \mathbf{r}')^2}{r^4} - \dots \right]$$

\leftarrow $O(r'/r)$
 \leftarrow $O((r'/r)^2)$

the other terms are higher order than $(r'/r)^2$

$$\frac{1}{|\mathbf{r} - \mathbf{r}'|} = \frac{1}{r} \left[1 - \frac{\mathbf{r} \cdot \mathbf{r}'}{r^2} + \frac{3}{2} \frac{(\mathbf{r} \cdot \mathbf{r}')^2}{r^4} - \dots \right]$$

\leftarrow $O(r'/r)$
 \leftarrow $O((r'/r)^2)$ and higher

$$\begin{aligned}
 \Phi(\vec{r}) &= \int \frac{\rho(\vec{r}')}{|\vec{r}-\vec{r}'|} d\tau' \\
 &= \int \frac{\rho(\vec{r}')}{r'} \left[1 - \frac{\vec{r} \cdot \vec{r}'}{r'^2} + \frac{1}{2} \frac{(\vec{r} \cdot \vec{r}')^2}{r'^4} - \dots \right] d\tau' \\
 &= \int \frac{\rho(\vec{r}')}{r'} d\tau' + \int \frac{\rho(\vec{r}')}{r'^2} \vec{r} \cdot \vec{r}' d\tau' + \int \frac{\rho(\vec{r}')}{r'^3} \left(\frac{3}{2} \frac{(\vec{r} \cdot \vec{r}')^2}{r'^2} - \frac{1}{2} r'^2 \right) d\tau' + \dots \\
 &= Q/r + \dots
 \end{aligned}$$

$$\Phi(\vec{r}) = \underbrace{Q/r}_{\text{monopole potential}} + \underbrace{\dots}_{\text{dipole potential}} + \underbrace{\dots}_{\text{quadrupole potential}} + \dots$$

This is the neatest form, because using ∇^2 exposes the powers of $1/r$, but it is easier to take the gradient of the 1st form.

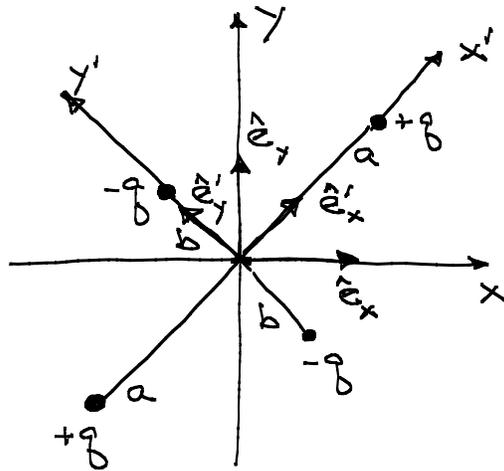
(c) Use $\nabla \left(\frac{1}{r} \right) = -\frac{\vec{r}}{r^3}$ and $\nabla x_i = \hat{e}_i$

$$\begin{aligned}
 \nabla \Phi &= \nabla \left(\frac{Q}{r} \right) + \nabla \left(\frac{p_i x_i}{r^2} \right) + \nabla \left(\frac{Q_{ij} x_i x_j}{r^3} \right) + \dots \\
 &= -\frac{Q \vec{r}}{r^3} + \dots
 \end{aligned}$$

Notice that $Q_{jk} = \sum_i Q_i x_i^j x_i^k = \sum_i Q_i x_i^j x_i^k = \sum_i Q_i x_i^k x_i^j = Q_{kj}$

$$= -\frac{Q}{r_1^3} + \frac{3Q}{r_1^5} \langle x_1^2 \rangle - \frac{3Q}{r_1^5} \langle x_1^2 \rangle + \dots$$

$$\langle x^2 \rangle = \frac{1}{Q} \sum_i Q_i x_i^2 = \frac{Q}{Q} + \frac{3Q \langle x_1^2 \rangle}{Q} - \frac{3Q \langle x_1^2 \rangle}{Q} + \dots$$



Total charge is zero
Dipole moment is like the center of mass and thus is zero.

(d) $Q_{jk} = \sum_{\alpha} q_{\alpha} [\omega_{\alpha} x_{\alpha}^j x_{\alpha}^k - \delta_{jk} (r_{\alpha})^2]$ α labels which charge

$$\sum_{\alpha} q_{\alpha} (r_{\alpha})^2 = \sum_{\alpha} q_{\alpha} a^2 - \sum_{\alpha} q_{\alpha} b^2 = \sum_{\alpha} q_{\alpha} (a^2 - b^2)$$

$$\sum_{\alpha} q_{\alpha} x_{\alpha}^2 x_{\alpha}^2 = \sum_{\alpha} q_{\alpha} a^4 - \sum_{\alpha} q_{\alpha} b^4 = \sum_{\alpha} q_{\alpha} (a^4 - b^4)$$

$$\sum_{\alpha} q_{\alpha} x_{\alpha}^2 y_{\alpha}^2 = \sum_{\alpha} q_{\alpha} a^2 b^2 - \sum_{\alpha} q_{\alpha} (-b^2 a^2) = \sum_{\alpha} q_{\alpha} (a^2 + b^2)$$

$$Q_{xx} = \sum_{\alpha} q_{\alpha} (a^4 - b^4) - \sum_{\alpha} q_{\alpha} (a^2 - b^2)^2 = \sum_{\alpha} q_{\alpha} (a^4 - b^4)$$

$$Q_{yy} = \sum_{\alpha} q_{\alpha} (a^4 - b^4)$$

$$Q_{xy} = Q_{yx} = \sum_{\alpha} q_{\alpha} (a^2 + b^2)$$

$$Q_{zz} = -\sum_{\alpha} q_{\alpha} (a^2 - b^2)$$

$$Q_{xz} = Q_{zx} = Q_{yz} = Q_{zy} = 0$$

Symmetric and traceless

(e) The matrix for \vec{Q} is

$$\begin{pmatrix} Q_{xx} & Q_{xy} & Q_{xz} \\ Q_{yx} & Q_{yy} & Q_{yz} \\ Q_{zx} & Q_{zy} & Q_{zz} \end{pmatrix} = \rho \begin{pmatrix} a^2 - b^2 & 2(a^2 + b^2) & 0 \\ 2(a^2 + b^2) & a^2 - b^2 & 0 \\ 0 & 0 & -2(a^2 - b^2) \end{pmatrix}$$

It is obvious that $\hat{e}_z = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$ is an eigenvector with eigenvalue $-2\rho(a^2 - b^2)$. Moreover, the symmetry of the charge distribution means that the other two eigenvectors are

$$\hat{e}'_x = \frac{1}{\sqrt{2}}(\hat{e}_x + \hat{e}_y) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix}$$

$$\hat{e}'_y = \frac{1}{\sqrt{2}}(-\hat{e}_x + \hat{e}_y) = \frac{1}{\sqrt{2}} \begin{pmatrix} -1 \\ 1 \\ 0 \end{pmatrix}$$

We get the eigenvalues from

$$\rho \begin{pmatrix} a^2 - b^2 & 2(a^2 + b^2) & 0 \\ 2(a^2 + b^2) & a^2 - b^2 & 0 \\ 0 & 0 & -2(a^2 - b^2) \end{pmatrix} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} = \underbrace{2\rho(a^2 + b^2)}_{\text{eigenvalue for } \hat{e}'_x} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix}$$

$$\rho \begin{pmatrix} a^2 - b^2 & 2(a^2 + b^2) & 0 \\ 2(a^2 + b^2) & a^2 - b^2 & 0 \\ 0 & 0 & -2(a^2 - b^2) \end{pmatrix} \begin{pmatrix} -1 \\ 1 \\ 0 \end{pmatrix} = \underbrace{-2\rho(a^2 - b^2)}_{\text{eigenvalue for } \hat{e}'_y} \begin{pmatrix} -1 \\ 1 \\ 0 \end{pmatrix}$$

(f) Now, using the primed coordinates,

$$Q'_{jk} = \sum_{\alpha} \rho_{\alpha} \left[\partial_{x_j} x_{\alpha}^{\prime} \partial_{x_k} x_{\alpha}^{\prime} - \delta_{jk} (r_{\alpha}^{\prime})^2 \right]$$

$$\sum_{\alpha} \rho_{\alpha} (r_{\alpha}^{\prime})^2 = 2\rho(a^2 - b^2) \quad \text{as before}$$

$$\sum_{\alpha} q_{\alpha} x_{\alpha}^{\prime} x_{\alpha}^{\prime} = \sum_{\alpha} q_{\alpha} a^2$$

Only the positive charges are on the x' axis

$$\sum_{\alpha} q_{\alpha} y_{\alpha}^{\prime} y_{\alpha}^{\prime} = -\sum_{\alpha} q_{\alpha} b^2$$

Only the negative charges are on the y' axis

$$\sum_{\alpha} q_{\alpha} x_{\alpha}^{\prime} y_{\alpha}^{\prime} = 0$$

no charge has both nonzero x' and y' coordinates

So

$$Q'_{xx} = 6q a^2 - 2q(a^2 - b^2) = 2q(2a^2 + b^2)$$

$$Q'_{yy} = -6q b^2 - 2q(a^2 - b^2) = -2q(a^2 + 2b^2)$$

$$Q'_{zz} = -2q(a^2 - b^2)$$

All off-diagonals are zero

