

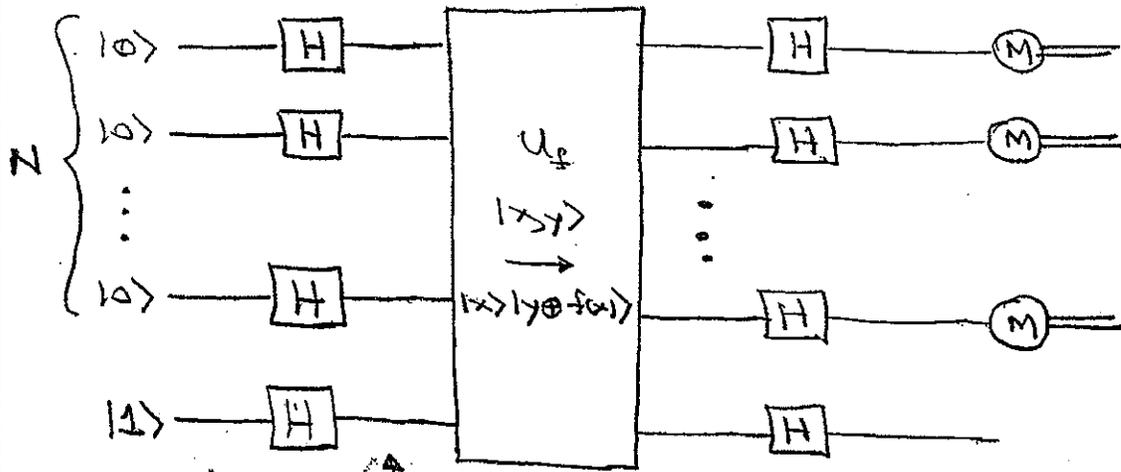
Quantum Computation

Lectures 13-14

Quantum Fourier transform and phase estimation

N-bit Boolean function  $f: \{0,1\}^N \rightarrow \{0,1\}$   
 Deutsch-Jozsa algorithm: which is either constant or balanced.

Number of classical calls to be sure is  $\sim \frac{2^N}{2} = 2^{N-1}$



$$|\psi_0\rangle = |0\rangle^{\otimes N} |1\rangle$$

$$|\psi_1\rangle = H^{\otimes N} |0\rangle^{\otimes N} \otimes H |1\rangle$$

$$= \frac{1}{\sqrt{2^N}} \sum_x |x\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

$$|\psi_2\rangle = \frac{1}{\sqrt{2^N}} \sum_x (-1)^{f(x)} |x\rangle \otimes \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

$f$  constant:  $\langle x | x \rangle = \pm |0\rangle^{\otimes N}$   
 $f$  balanced:  
 $\langle 0 | x \rangle = \frac{1}{\sqrt{2^N}} \sum_x (-1)^{f(x)} = 0$

One-shot determination of whether function is constant or balanced.

$$|\psi_2\rangle = \frac{1}{\sqrt{2^N}} \sum_x |x\rangle \frac{1}{\sqrt{2}} (|f(x)\rangle - |\neg f(x)\rangle)$$

quantum parallelism

$$= \frac{1}{\sqrt{2^N}} \sum_x (-1)^{f(x)} |x\rangle \otimes \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

ancilla qubit left unentangled.

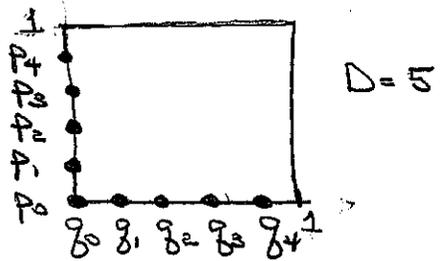
phase kickback

function value written in phase

$$\langle \phi^{\text{constant}} | \phi^{\text{balanced}} \rangle = \frac{1}{\sqrt{2^N}} \sum_x (-1)^{f(x)} = 0$$

# Quantum Fourier transform

D-dimensional Hilbert space



Orthonormal "position" basis

$$|q_j\rangle = |e_j\rangle = |j\rangle, \quad j = 0, \dots, D-1$$

$$q_j = j/D$$

use this notation when there's no other basis around to cause confusion

$$\text{with } D = (\text{area of phase space}) = 1$$

Conjugate "momentum" basis

$$|p_k\rangle, \quad k = 0, \dots, D-1$$

$$p_k = k/D$$

j and k are defined mod D; the FT coefficient does mod D arithmetic.

$$|p_k\rangle = \frac{1}{\sqrt{D}} \sum_{j=0}^{D-1} |q_j\rangle e^{2\pi i j k / D}$$

$$\langle q_j | p_k \rangle = \frac{1}{\sqrt{D}} \exp\left(\frac{i}{\hbar} q_j p_k\right) = \frac{1}{\sqrt{D}} e^{2\pi i j k / D}$$

$$\langle q_j | p_k \rangle = \langle q_k | p_j \rangle = \langle q_j | p_k \rangle^* = \langle p_k | q_j \rangle$$

discrete FT coefficients

periodic in j (q) with period  $\Delta j = D/k = 1/p_k$   
 $(\Delta q = 1/k = 1/D p_k = 2\pi \hbar / p_k)$

Matrix elements of a unitary matrix:

$$\frac{1}{D} \sum_{l=0}^{D-1} e^{2\pi i l (j-k) / D} = \frac{1}{D} \frac{1 - e^{2\pi i (j-k)}}{1 - e^{2\pi i (j-k) / D}} = \delta_{jk}$$

Quantum FT operator:  $F |q_j\rangle = |p_j\rangle$

$$\langle q_j | F | q_k \rangle = \langle q_j | p_k \rangle = \frac{1}{\sqrt{D}} e^{2\pi i j k / D}$$

$$y_k = \langle p_k | \psi \rangle = \langle q_k | F^\dagger | \psi \rangle = \frac{1}{\sqrt{D}} \sum_{j=0}^{D-1} e^{-2\pi i j k / D} \underbrace{\langle q_j | \psi \rangle}_{x_j}$$

$$y_k = \frac{1}{\sqrt{D}} \sum_{j=0}^{D-1} e^{-2\pi i j k / D} x_j \quad \text{Discrete FT}$$

$$F|P_j\rangle = \sum_k F|g_k\rangle \langle g_k|P_j\rangle = \sum_k |P_k\rangle \langle P_k|g_j\rangle = |g_j\rangle$$

$$\Rightarrow F^2|g_j\rangle = |g_j\rangle \Rightarrow F^2 = 1$$

↑  
parity operator

F has eigenvalues  
 $\pm 1, \pm i$

What does this have to do with qubits? N qubits  
 $D = 2^N$

Notation:  $j = j_1 \dots j_N = \sum_{x=1}^N j_x 2^{N-x}$

$$|g_j\rangle = |j\rangle = |j_1\rangle \otimes \dots \otimes |j_N\rangle$$

binary representation of  $j$  ( $j_x = 0, 1$ )

$$g_j = j/D = j/2^N = 0.j_1 \dots j_N = \sum_{x=1}^N j_x 2^{-x}$$

$$j/2^k = j_1 \dots j_{N-k} \cdot j_{N-k+1} \dots j_N = \sum_{x=1}^N j_x 2^{N-x-k}$$

$$e^{2\pi i j/2^k} = e^{2\pi i (0.j_{N-k+1} \dots j_N)}$$

### Key qubit representation of F:

$$\begin{aligned}
 |P_j\rangle &= \frac{F}{\sqrt{N}} |g_j\rangle = \frac{F}{\sqrt{N}} |j\rangle = \frac{1}{\sqrt{N/2}} \sum_{k=0}^{2^N-1} |k\rangle e^{2\pi i k j / D} \\
 &\quad \uparrow \\
 &\text{FT on N qubits} \\
 &= \frac{1}{\sqrt{N/2}} \sum_{k_1, \dots, k_N=0}^1 |k_1\rangle \otimes \dots \otimes |k_N\rangle \exp \left[ 2\pi i \left( \sum_{x=1}^N k_x 2^{N-x} \right) \frac{j}{D} \right] \\
 &\quad \downarrow \\
 &= \frac{1}{\sqrt{N/2}} \bigotimes_{x=1}^N \sum_{k_x=0}^1 |k_x\rangle e^{2\pi i k_x j / 2^x} \\
 &= e^{2\pi i (0.j_{N-x+1} \dots j_N) k_x}
 \end{aligned}$$

$$F_N |j\rangle = \frac{1}{\sqrt{2^N}} \bigotimes_{l=1}^N (|0\rangle + e^{2\pi i j / 2^l} |1\rangle)$$

↑  
product state

$$F_N |j\rangle = \frac{1}{\sqrt{2}} (|0\rangle + e^{2\pi i j / 2^1} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + e^{2\pi i j / 2^2} |1\rangle) \otimes \dots \otimes \frac{1}{\sqrt{2}} (|0\rangle + e^{2\pi i j / 2^N} |1\rangle)$$

$$= \frac{1}{\sqrt{2}} (|0\rangle + e^{2\pi i (0 \cdot j_N)} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + e^{2\pi i (0 \cdot j_{N-1} + j_N)} |1\rangle) \otimes \dots \otimes \frac{1}{\sqrt{2}} (|0\rangle + e^{2\pi i (0 \cdot j_1 + \dots + j_N)} |1\rangle)$$

This form is crucial because it shows that  $F$  can be implemented by separate (controlled) operations on each qubit, giving an  $O(N^2)$  algorithm.

Building circuits for  $F_N$ :

Define  $R_k \equiv e^{2\pi i / 2^{k+1}} \underbrace{e^{-iZ 2\pi / 2^{k+1}}}_{\text{rotation by } 2\pi / 2^k = \pi / 2^{k-1} \text{ about } \vec{e}_z}$   $\longleftrightarrow$   $\begin{pmatrix} 1 & 0 \\ 0 & e^{+2\pi i / 2^k} \end{pmatrix}$

$R_0 = I$   
 $R_1 = Z$   
 $R_2 = S$   
 $R_3 = T$   
...

$$H|a\rangle = \frac{1}{\sqrt{2}} (|0\rangle + (-1)^a |1\rangle)$$

$$= \frac{1}{\sqrt{2}} (|0\rangle + e^{i\pi a} |1\rangle)$$

$$= \frac{1}{\sqrt{2}} (|0\rangle + e^{2\pi i (0.a)} |1\rangle)$$

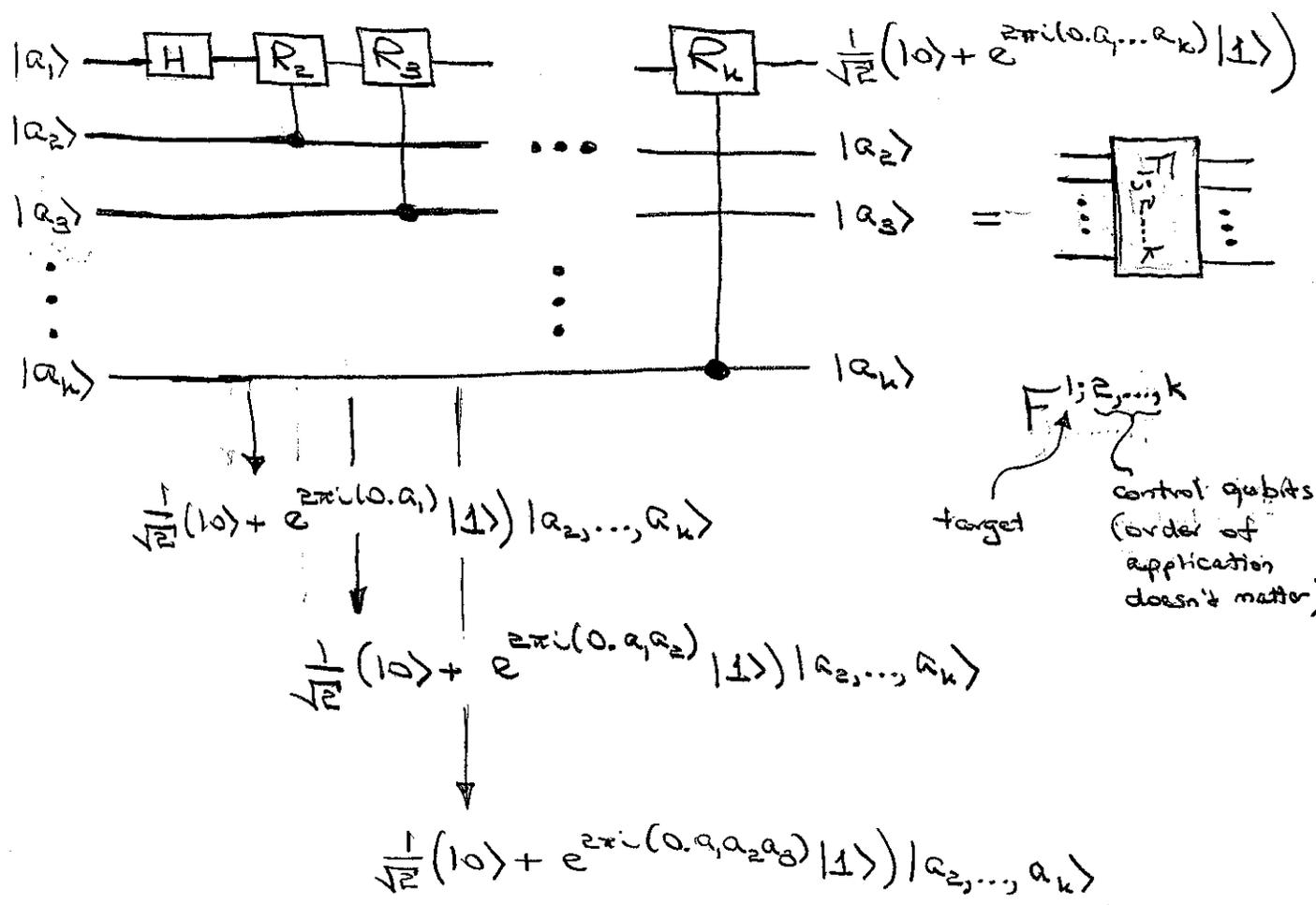
$$R_k^a (\alpha |0\rangle + \beta |1\rangle) = \alpha |0\rangle + \beta e^{2\pi i a / 2^k} |1\rangle$$

$$= \alpha |0\rangle + \beta e^{2\pi i (0.\underbrace{0 \dots 0}_k a)} |1\rangle$$

k-1 0's

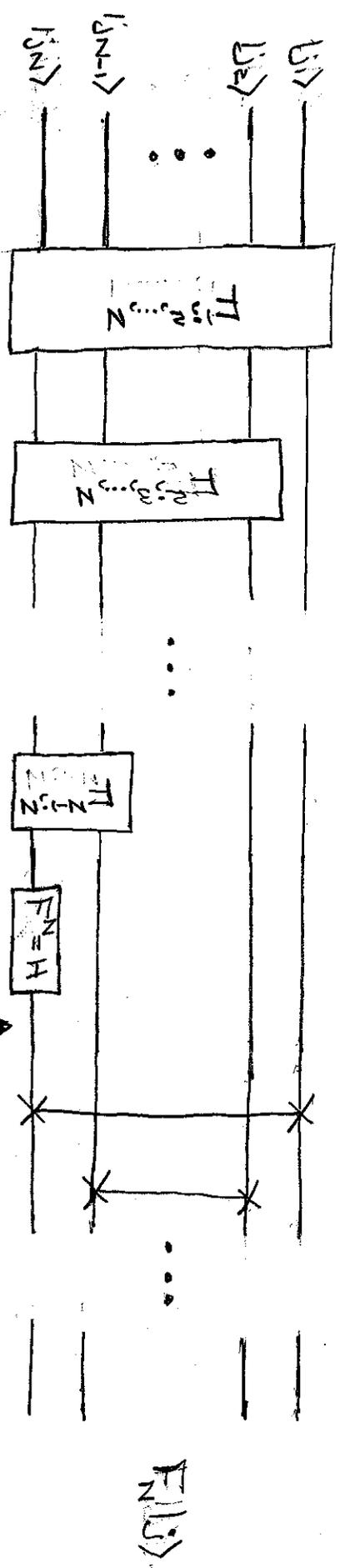
$$R_{k+1}^2 = R_k$$

The basic building block of the circuit is



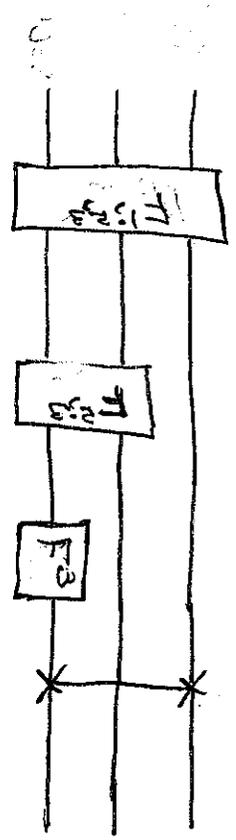
Resources for  $F_{|a_2, \dots, a_k}$  :  $O(k)$

FT Circuits:



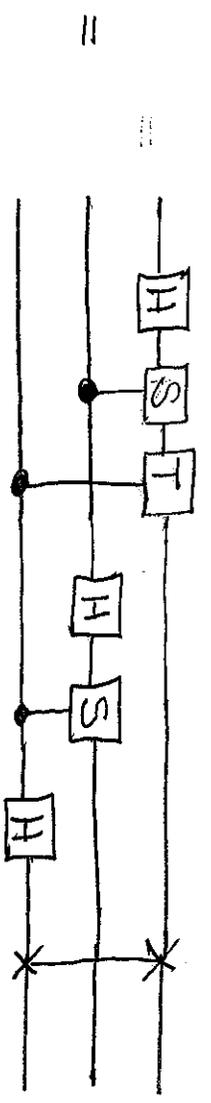
$$\frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i(0 \cdot j_1 - j_1 N)} |1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i(0 \cdot j_2 - j_2 N)} |1\rangle) \otimes \dots \otimes \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i(0 \cdot j_{N-1} - j_{N-1} N)} |1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i(0 \cdot j_N - j_N N)} |1\rangle)$$

N=3



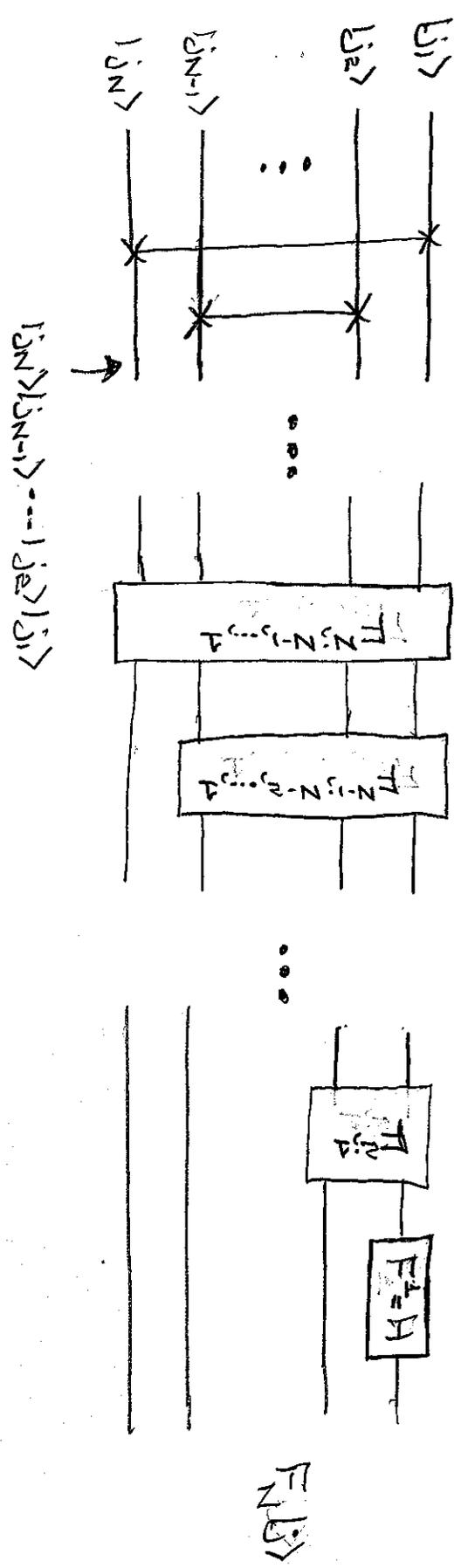
You can have the Swap gates out by swapping the labels of the qubits instead

Resources:  $O(N^2)$

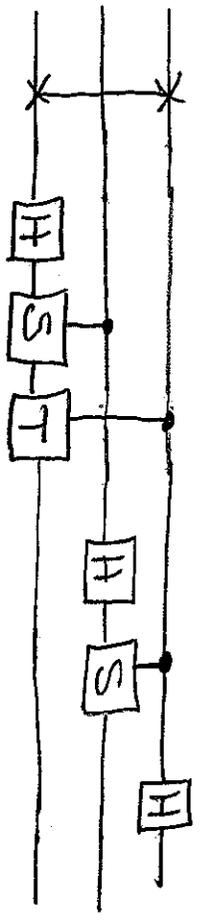


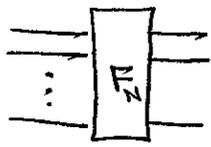
①

② You can also put the swaps first, which you can see either directly or by moving the swaps through the  $F^i$ 's.



$N=3$



Now we can simply write  when we want to include a Fourier transform.

Phase estimation

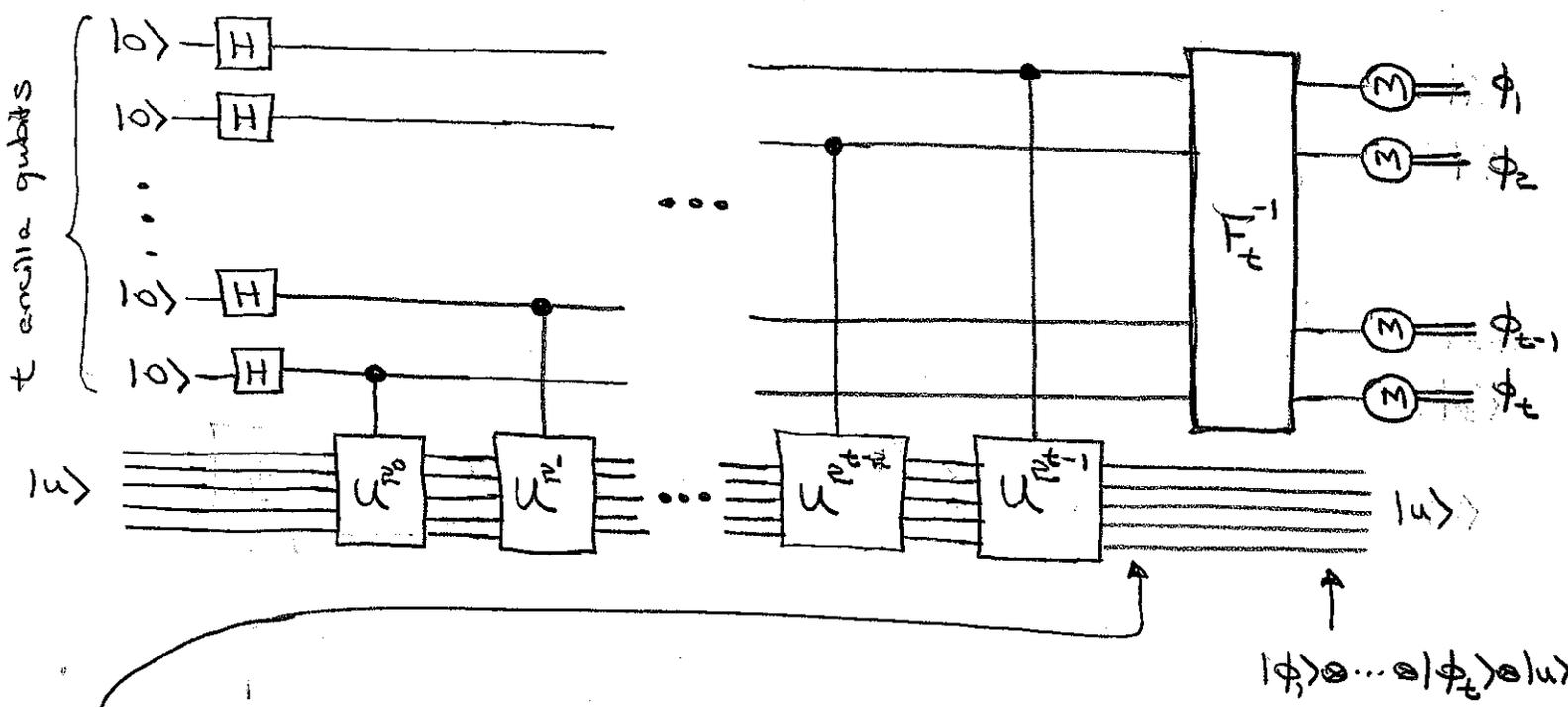
unknown phase as fraction of  $2\pi$

$$U|u\rangle = e^{2\pi i \phi} |u\rangle$$

↑  
 unitary such that  $U^{2^l}$ ,  $l=0, \dots, t-1$ , can be implemented efficiently.  
 ↳ eigenvector of  $U$  that can be prepared reliably

① Assume that  $\phi = 0.\phi_1 \dots \phi_t = \phi_1 \dots \phi_t / 2^t$

Phase estimation circuit



$$\begin{aligned}
 &= \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i (2^{t-1} \phi)} |1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i (2^{t-2} \phi)} |1\rangle) \otimes \dots \otimes \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i (2^1 \phi)} |1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i (2^0 \phi)} |1\rangle) \otimes |u\rangle \\
 &= \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i (0.\phi_t)} |1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i (0.\phi_{t-1}\phi_t)} |1\rangle) \otimes \dots \otimes \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i (0.\phi_2 \dots \phi_t)} |1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i (0.\phi_1 \dots \phi_t)} |1\rangle) \otimes |u\rangle \\
 &= F_t |\phi\rangle = F_t | \phi_1 \rangle \otimes | \phi_2 \rangle \otimes \dots \otimes | \phi_t \rangle
 \end{aligned}$$

The controlled unitaries prepare a momentum state

$$|P_{z^t, \phi}\rangle = F_z |\phi\rangle = \frac{1}{\sqrt{D}} \sum_{k=0}^{z^t-1} e^{2\pi i k \phi} |k\rangle$$

periodic in position with period  $\Delta q = 1/\phi = 2\pi\hbar/P_z = \hbar/P_z$

We can determine  $\phi$ , i.e., the period  $2\pi\hbar/P_z$ , from a measurement in the momentum basis, but we don't assume we know how to do that. So we do a FT to put the phase information in the standard basis (the inverse FT is an interferometer with  $z^t$  inputs and outputs).

② What happens when  $\phi = .\phi_1\phi_2\dots$  has more than  $t$  digits?  
 Now let  $|\phi\rangle$  denote the state that is input to the inverse FT:

$$|\phi\rangle = \frac{1}{\sqrt{z^t}} \bigotimes_{l=1}^t (|0\rangle + e^{2\pi i (z^{t-l} \phi)} |1\rangle)$$

Reverse the steps leading to the qubit form of the FT:

$$|\phi\rangle = \frac{1}{\sqrt{z^t}} \bigotimes_{l=1}^t (|0\rangle + e^{2\pi i (z^{t-l} \phi)/z^l} |1\rangle)$$

$$\underbrace{\sum_{k=0}^{z^l-1} |k\rangle e^{2\pi i k (z^{t-l} \phi)/z^l}}_j$$

$$= \frac{1}{\sqrt{z^t}} \sum_{k_1, \dots, k_N} |k_1\rangle \otimes \dots \otimes |k_N\rangle \exp \left[ 2\pi i \underbrace{\left( \sum_{l=1}^t k_l z^{t-l} \right)}_k \phi \right]$$

$$= \frac{1}{\sqrt{z^t}} \sum_{k=0}^{z^t-1} |k\rangle e^{2\pi i k \phi}$$

$$|0\rangle^{\otimes t} |u\rangle \rightarrow |\phi\rangle |u\rangle = \frac{1}{\sqrt{z^t}} \sum_{k=0}^{z^t-1} |k\rangle U^k |u\rangle$$

← This is the state before the FT

$$\langle q_j | F^\dagger \phi \rangle = \frac{1}{\sqrt{z^t}} \sum_{k=0}^{z^t-1} e^{2\pi i k (\phi - j/z^t)}$$

The calculation at the bottom of page 9 shows that

$$|0\rangle^{\otimes t} |u\rangle \rightarrow \frac{1}{\sqrt{2^{t/2}}} \sum_{k=0}^{2^t-1} |k\rangle U^k |u\rangle,$$

which implies

$$|0\rangle^{\otimes t} |\psi\rangle \rightarrow \frac{1}{\sqrt{2^{t/2}}} \sum_{k=0}^{2^t-1} |k\rangle U^k |\psi\rangle.$$

We can get at this result more directly by working with the circuit. After the Hadamards, the controlled unitaries do the transformation

$$|j\rangle |\psi\rangle \rightarrow |j\rangle \otimes \dots \otimes |j_t\rangle \underbrace{U^{2^{t-1}j_1} \dots U^{2^{j_{t-1}}} U^{2^{j_t}}}_{U^{\sum_{k=1}^t 2^{t-k} j_k} = U^j} |\psi\rangle$$

$$|j\rangle |\psi\rangle \rightarrow |j\rangle U^j |\psi\rangle$$

So with the Hadamards, the transformation is

$$|0\rangle^{\otimes t} |\psi\rangle \rightarrow \frac{1}{\sqrt{2^{t/2}}} \sum_{j=0}^{2^t-1} |j\rangle U^j |\psi\rangle.$$

So, if  $|\psi\rangle$  is a superposition of several eigenstates of  $U$ , the output measurement will yield one of the eigenvalues with the probability from the superposition.

$$\langle g_j | F^t | \phi \rangle = \frac{1}{2^t} \sum_{k=0}^{2^t-1} (e^{2\pi i(\phi - j)/2^t})^k$$

$$= \frac{1}{2^t} \frac{1 - e^{2\pi i(\phi - j)}}{1 - e^{2\pi i(\phi - j)/2^t}}$$

$b \equiv \lfloor \phi 2^t \rfloor$ , and  $\phi 2^t = b + \delta$ ,  $0 \leq \delta < 1$ ,  $0 \leq \delta < 1$

$\phi = b 2^{-t} + \delta 2^{-t}$   
 ↑                      ↑  
 t-bit approximation to  $\phi$       error in the t-bit approximation

$$\langle g_{b+\ell} | F^t | \phi \rangle = \frac{1}{2^t} \frac{1 - e^{2\pi i(\delta - \ell)}}{1 - e^{2\pi i(\delta - \ell)/2^t}}$$

$$= \frac{1}{2^t} \frac{e^{i\pi(\delta - \ell)/2} \sin(\pi(\delta - \ell)/2^t)}{e^{i\pi(\delta - \ell)/2^t} \sin(\pi(\delta - \ell)/2^t)}$$

$$| \langle g_{b+\ell} | F^t | \phi \rangle |^2 = \frac{1}{2^{2t}} \frac{\sin^2 \pi(\delta - \ell)}{\sin^2 \pi(\delta - \ell)/2^t}$$

$\delta 2^{-t}$  is error in determination of  $\phi$  because  $\phi$  has more than t bits.  
 $\ell 2^{-t}$  is error in determination of  $\phi$  because the measurement doesn't yield  $b$ .

$| \ell | \leq 2^{n-1}$  means an error  $\leq n$  bits, giving  $\phi$  to  $N = t - n$  bits.

$$P(| \ell | > 2^{n-1}) = \sum_{\ell = -2^{n-1}+1}^{-2^{n-1}-1} | \langle g_{b+\ell} | F^t | \phi \rangle |^2 + \sum_{\ell = 2^{n-1}+1}^{2^t-1} | \langle g_{b+\ell} | F^t | \phi \rangle |^2$$

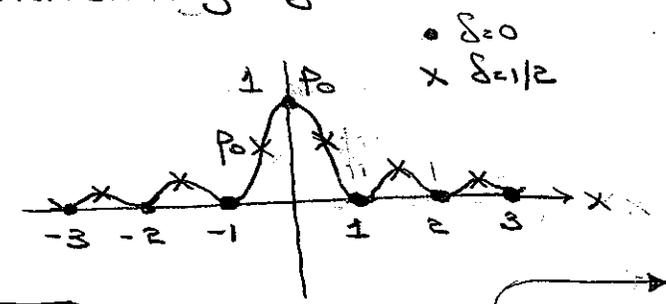
(probability of getting less than  $N = t - n$  bits)

$\ell = -2^{n-1}+1, \dots, 2^{n-1}$

(center summing range at  $\ell=0$ )

$$| \langle g_{b+\ell} | F^t | \phi \rangle |^2 = \frac{1}{2^{2t}} \frac{\sin^2 \pi x}{\sin^2 \pi x / 2^t}$$

$x = \delta - \ell$

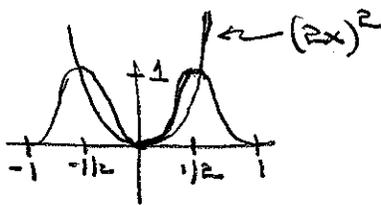


Bound below is based on worst case ( $\delta = 1/2$ ).

recurrences

$$\sin^2 \pi x \leq 1$$

$$\sin^2 \pi x \geq (Rx)^2$$



$$\frac{l-\delta}{2^t} \leq \frac{1}{2^t} (2^{t-1} - \delta) = \frac{1}{2} - \frac{\delta}{2^t} \leq \frac{1}{2}$$

$$\frac{l-\delta}{2^t} \geq \frac{1}{2^t} (-2^{t-1} + 1 - \delta) = -\frac{1}{2} + \frac{1-\delta}{2^t} \geq -\frac{1}{2}$$

$$\left| \frac{l-\delta}{2^t} \right| \leq \frac{1}{2}$$

$$\Rightarrow \sin^2 \left( \frac{\pi(l-\delta)}{2^t} \right) \geq \left( \frac{2(l-\delta)}{2^t} \right)^2 = \frac{4(l-\delta)^2}{2^{2t}}$$

$$| \langle \psi_{b+l} | F^t | \phi \rangle |^2 \leq \frac{1}{4(l-\delta)^2}$$

$$P(12) > 2^{2t-1} \leq \frac{1}{4} \left( \sum_{l=-2^{t-1}+1}^{-2^{t-1}-1} \frac{1}{(l-\delta)^2} + \sum_{l=2^{t-1}+1}^{2^{t-1}-1} \frac{1}{(l-\delta)^2} \right)$$

$$\leq \frac{1}{2^2} \leq \frac{1}{(l-1)^2}$$

$$\leq \frac{1}{4} \left( \sum_{l=-2^{t-1}+1}^{-2^{t-1}-1} \frac{1}{l^2} + \sum_{l=2^{t-1}+1}^{2^{t-1}-1} \frac{1}{l^2} \right)$$

$$= \frac{1}{4} \sum_{l=2^{t-1}}^{2^{t-1}-1} \frac{1}{l^2}$$

$$\leq \frac{1}{4} \int_{2^{t-1}}^{2^{t-1}-1} \frac{dl}{l^2}$$

$$\leq \frac{1}{4} \int_{2^{t-1}}^{\infty} \frac{dl}{l^2}$$

$$= \frac{1}{4} \left[ \frac{1}{2^{t-1}-1} - \frac{1}{2^{t-1}} \right]$$

$$E \equiv \left( \begin{array}{l} \text{probability of getting less} \\ \text{than } N = t-n \text{ bits of } \phi \end{array} \right) \leq \frac{1}{R} \frac{1}{R^{t-n-1}} = \frac{1}{R} \frac{1}{R^{t-N-1}}$$

$$= 1 - \left( \begin{array}{l} \text{probability of getting} \\ N = t-n \text{ bits or more} \end{array} \right)$$

$$n \leq \log(R + E^{-1})$$