

Reminder: Wiki Tutorial 7pm tomorrow  
 will commit <sup>tips</sup> suggestions - select topic by next Thu.  
 HW assignment - next two days

①

## Ph 452 Lecture 21

### 1Q1 Quick Review

- Error channels: Bit-flip, phase-flip = phase damping
- Amplitude-damping channel

$$A_0 = \begin{bmatrix} 1 & 0 \\ 0 & \sqrt{1-p} \end{bmatrix}, A_1 = \begin{bmatrix} 0 & \sqrt{p} \\ 0 & 0 \end{bmatrix}$$

Bloch sphere:  $\frac{1}{2}(\mathbf{I} + \vec{r} \cdot \vec{\sigma})$

$$\vec{r} \rightarrow (\sqrt{1-p} r_x, \sqrt{1-p} r_y, (1-p)r_z + \underline{\underline{p}})$$

$$A_0 Z A_0^\dagger + A_1 Z A_1^\dagger = (1-p)Z, \text{ but also}$$

$$A_0 \mathbf{I} A_0^\dagger + A_1 \mathbf{I} A_1^\dagger = \begin{pmatrix} 1+p & 0 \\ 0 & 1-p \end{pmatrix} = \mathbf{I} + \underline{\underline{pZ}} \quad \text{non-unitary}$$

### Depolarizing channel:

$$\rho \rightarrow (1-p)\rho + \frac{p}{3}(X\rho X + Y\rho Y + Z\rho Z) \quad 0 \leq p \leq 1$$

Factoid:  $\frac{1}{3}(\mathbf{I}\rho\mathbf{I} + X\rho X + Y\rho Y + Z\rho Z) = \mathbf{I}$

Proof: Use  $\rho = \frac{1}{2}(\mathbf{I} + \vec{r} \cdot \vec{\sigma})$

$$\Rightarrow \rho \rightarrow (1-p)\rho + \frac{p}{3}(2\mathbf{I} - \rho) = (1 - \frac{4}{3}p)\rho + \frac{4}{3}p(\frac{\mathbf{I}}{2})$$

$$\Rightarrow \text{with "probability" } (1 - \frac{4}{3}p) \text{ nothing happens, "probab"} \frac{4}{3}p \rho \rightarrow \mathbf{I}/2, \quad (p \leq 3/4)$$

(note:  $\mathbf{I}$  contains a component of  $\rho$ , so elements of "ensemble" are not independent.)

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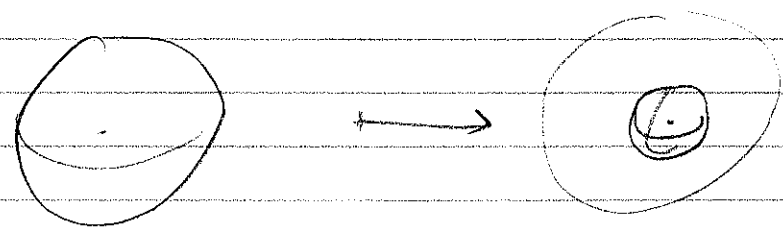
Bloch sphere picture

$$\frac{1}{2}(\mathbb{I} + \vec{r} \cdot \vec{\sigma}) \rightarrow (1 - \frac{4}{3}\rho) \left(\frac{1}{2}\right) (\mathbb{I} + \vec{r} \cdot \vec{\sigma}) + \frac{4}{3}\rho \left(\frac{1}{2}\right) \mathbb{I}$$

$$= \frac{1}{2} \mathbb{I} + (1 - \frac{4}{3}\rho) \vec{r} \cdot \vec{\sigma}$$

Q: If  $\rho > 3/4$ ,  
 does Bloch ball  
 "invert"?  
 A: No.  
 "can fit inside  
 class"

$$\vec{r} \rightarrow (1 - \frac{4}{3}\rho) \vec{r}$$



Bloch ball  
 uniformly  
 contracts.

Open Q. Systems

- ①  $|\psi\rangle \rightarrow \rho$
- ②  $U^\dagger U = \mathbb{I} \rightarrow \sum A_i^\dagger A_i = \mathbb{I}$      Aside: Davies Thm;  $d^2$  Kraus ops suffice
- ③  $\sum \Pi_i = \mathbb{I}, \Pi_i \Pi_j = \delta_{ij}, \Pi_j, \Pi_i = \Pi_j^\dagger \rightarrow ?$

Q: How big does  
 system B have to  
 be?  
 A:  $d^2$   
 (From Davies Thm?)

$$\rho_A \rightarrow \frac{1}{\text{prob}(K)} \text{Tr}_B \left[ (\mathbb{I}_A \otimes \Pi_K^{(B)}) U_{AB} (\rho_A \otimes |\psi\rangle_B \langle \psi|) U_{AB}^\dagger (\mathbb{I}_A \otimes \Pi_K^{(B)}) \right]$$

$$A_{jK} = \langle 0 | \Pi_K^{(B)} U_{AB} | \psi \rangle_B$$

$$\rho_A \rightarrow \sum_j A_{jK} \rho_A A_{jK}^\dagger / \text{prob}(K)$$

$$\sum_{JK} A_{jK}^\dagger A_{jK} = \sum_{JK} \langle \psi | U_{AB}^\dagger \Pi_K^{(B)} | \psi \rangle_B \langle 0 | \Pi_K^{(B)} U_{AB} | \psi \rangle_B$$

$$= \mathbb{I}_A$$

$$\text{prob}(K) = \text{Tr}_B \left[ (\mathbb{I}_A \otimes \Pi_K) U_{AB} (\rho_A \otimes |\psi\rangle \langle \psi|) U_{AB}^\dagger (\mathbb{I}_A \otimes \Pi_K) \right]$$

$$= \text{Tr}_A \left( \sum_j A_{jK}^\dagger A_{jK} \rho_A \right)$$

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$$\text{Let } E_k := \sum_j A_{jk}^\dagger A_{jk}$$

$$\text{Then } \boxed{\sum_k E_k = \mathbb{I}, \text{ Prob}(k) = \text{tr}(E_k \rho), E_k \geq 0}$$

↑  
"measure"
↑  
Positive operator

POVM: Positive operator valued measure!

Example: One-qubit POVM



$$E_k = \frac{1}{3} (\mathbb{I} + \hat{n}_k \cdot \vec{\sigma})$$

$$\sum E_k = 3 \left( \frac{\mathbb{I}}{3} \right) + \left( \sum \hat{n}_k \right) \cdot \vec{\sigma} = \mathbb{I}, \quad E_k \geq 0$$

Note:  $\rho$  does change after a POVM, but the POVM formalism doesn't give a unique specification for the changes.

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POVM example:

Bob has  $|\psi_1\rangle = |0\rangle$  or  $|\psi_2\rangle = |1\rangle$ . Which does he have?  
 $= \frac{1}{\sqrt{2}}(|\psi_1\rangle + |\psi_1^\perp\rangle)$

Any projective meas. will mis-identify sometimes

Proof:

Let  $\{E_k\}$  be a projective meas. where outcome  $k$  implies Bob guesses " $|\psi_k\rangle$ ."

Then b/c  $|\psi_2\rangle$  has a  $|\psi_1\rangle$  component, Bob will sometimes guess " $|\psi_1\rangle$ " when given  $|\psi_2\rangle$

Now consider POVM:

$$E_1 = \frac{\sqrt{2}}{1+\sqrt{2}} |1\rangle\langle 1|, E_2 = \frac{\sqrt{2}}{1+\sqrt{2}} \frac{(|1\rangle\langle 0| + |0\rangle\langle 1|)}{2}, E_3 = I - E_1 - E_2$$

→ Distinguishes  $|\psi_1\rangle$  from  $|\psi_2\rangle$  some of the time, but never mis-identifies.

Q: What about  $|0\rangle, |1\rangle, |1\rangle$ , which determines  $|0\rangle$  some of the times, but never mis-identify

Proof: Given  $|\psi_1\rangle$ ,  $\text{prob}(1) = \text{tr}(E_1 |\psi_1\rangle\langle\psi_1|) = 0$

⇒ if outcome "1" obtained, must have been  $|\psi_2\rangle$

Given  $|\psi_2\rangle$ ,  $\text{prob}(2) = 0$  similarly

⇒ if outcome "2" obtained, must have been  $|\psi_1\rangle$

IF outcome "3" obtained → do decision.

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"Back to the main line of development: QEC"

"Generalized states, dynamics, measurements"

Quantum error:  $|\psi_0\rangle\langle\psi_0| \rightarrow \sum A_i |\psi_0\rangle\langle\psi_0| A_i^\dagger = \rho_{\text{final}}$

How bad is it, doc?

Purity:  $\text{tr } \rho^2 = \begin{cases} 1 & \rho = |\psi\rangle\langle\psi| \text{ for some } |\psi\rangle \\ \leq 0 & \text{otherwise} \end{cases}$  (recall  $\Pi_i^2 = \Pi_i$ )

Distance between qubits?  $d(\rho, \sigma)$ ?

Step 1: distance between pbits? (pbits)

"Trace distance" or "Kolmogorov distance"

$D(\vec{p}, \vec{q}) := \frac{1}{2} \sum_i |p_i - q_i|$

Features:  $D(p, q) = D(q, p)$ ,  $D(p, p) = 0$ ,  $D(p, q) + D(q, r) \geq D(p, r)$   
→ a "metric"

Motivation: Maximum difference in probabilities of an event occurring:

$D(\vec{p}, \vec{q}) = \max_s |p(s) - q(s)| = \max_s \left| \sum_{i \in S} p_i - \sum_{i \in S} q_i \right|$

Another measure: Fidelity:

$F(\vec{p}, \vec{q}) = \sum_i \sqrt{p_i q_i}$

Features: Not a metric,  $F(\vec{p}, \vec{p}) = 1$ ,  $F \geq 0$

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For quantum states:

Trace distance:  $D(\rho, \sigma) = \frac{1}{2} \text{Tr} |\rho - \sigma|$

recall  $|A| = \sqrt{A^\dagger A}$

For qubits:

$$D(\rho, \sigma) = \frac{1}{2} \text{tr} \left| \frac{1}{2} (\mathbb{I} + \vec{r} \cdot \vec{\sigma}) - \frac{1}{2} (\mathbb{I} + \vec{s} \cdot \vec{\sigma}) \right|$$
  
$$= \frac{1}{4} \text{tr} |(\vec{r} - \vec{s}) \cdot \vec{\sigma}|$$

$(\vec{r} - \vec{s}) \cdot \vec{\sigma}$  has eigen vals  $\pm |\vec{r} - \vec{s}|$ , so

$$D(\rho, \sigma) = \frac{1}{2} |\vec{r} - \vec{s}|$$

$\frac{1}{2}$  Euclidean dist btw  $\vec{r}, \vec{s}$ !

Got this far

Features: ①  $D$  is a metric, ②  $D(U\rho U^\dagger, U\sigma U^\dagger) = D(\rho, \sigma)$

③  $D = \max_{\text{POVMs } \{B_i\}} \text{tr} [E \rho - \sigma]$

"max diff. in probabilities any POVM will give"

④  $D(\mathcal{A}(\rho), \mathcal{A}(\sigma)) \leq D(\rho, \sigma)$

"contractive under quantum operations"

⑤ Convex:  $D(\sum p_i \rho_i, \sigma) \leq \sum p_i D(\rho_i, \sigma)$