

W.Ki, HW post: soon.

Errata: $F = \frac{\text{Tr}}{=} \sqrt{p^{1/2} \rho p^{1/2}}$, not $\sqrt{\text{tr} p^{1/2} \rho p^{1/2}}$
 $E_k \neq U_k E_k U_k^\dagger$

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PH 452 Lecture 22

1Q1 Quick Review

• Depolarizing channel

$$p = \left(1 - \frac{4}{3}p\right)p + \left(\frac{4}{3}p\right)\frac{I}{2} \quad 0 \leq p \leq 1$$

$\frac{I}{2}$ has a component of p in it, so $p, I/2$ are not independent.

$p > 3/4$ does not give an "inversion."

• POVM: $\sum E_k = I, E_k \geq 0, \text{prob}(k) = \text{tr}(E_k p)$

State after a POVM: wlog $E_k = A_k^\dagger A_k$, so

$$p \rightarrow \frac{A_k p A_k^\dagger}{\text{tr}(A_k^\dagger p A_k)}$$

[Not unique, $E_k = U_k^\dagger A_k^\dagger A_k U_k$
 $E_k = \sum_j A_{jk}^\dagger A_{jk}$]

New today →

Neumark's Dilation Theorem: All POVMs can be "lifted" to a projective measurement on a larger space

Davis' 1hr

Stinespring Dilation Theorem: All quantum operations can be "lifted" to unitary dynamics on a larger space

• POVM Application: Unambiguous State Discrimination

Input: $|\psi_1\rangle$ or $|\psi_2\rangle$ w/ 50% $\langle \psi_1 | \psi_2 \rangle \neq 0$

Object: Maximize probability that measurement is conclusive

① $\{|\psi_1\rangle\langle\psi_1|, |\psi_1\rangle\langle\psi_2| + |\psi_2\rangle\langle\psi_1|\}$ → $p = \frac{1}{2}(1 - |\langle\psi_1|\psi_2\rangle|^2) = \frac{1}{2} - \frac{1}{2}|\langle\psi_1|\psi_2\rangle|^2$

② $\{|\psi_1\rangle\langle\psi_2|, |\psi_2\rangle\langle\psi_1| + |\psi_1\rangle\langle\psi_2|\}$ → " " "

③ $E_1 = \frac{1 - |\psi_1\rangle\langle\psi_1|}{1 + |\psi_1\rangle\langle\psi_2|}, E_2 = \frac{1 - |\psi_2\rangle\langle\psi_2|}{1 + |\psi_2\rangle\langle\psi_1|}, E_3 = I - E_1 - E_2 \rightarrow p = 1 - |\langle\psi_1|\psi_2\rangle|^2$

New wish today →

$$\textcircled{3} E_1 = \frac{1 - |\psi_1 \rangle \langle \psi_1|}{1 + |\langle \psi_1 | \psi_2 \rangle|} \quad E_2 = \frac{1 - |\psi_2 \rangle \langle \psi_2|}{1 + |\langle \psi_1 | \psi_2 \rangle|}, \quad E_3 = 1 - E_1 - E_2$$

$$\begin{aligned} \text{Prob}(1 | \psi_2) &= \text{Tr}(E_1 |\psi_2 \rangle \langle \psi_2|) = \langle \psi_2 | E_1 | \psi_2 \rangle \\ &= \frac{1 - |\langle \psi_1 | \psi_2 \rangle|^2}{1 + |\langle \psi_1 | \psi_2 \rangle|} = 1 - |\langle \psi_1 | \psi_2 \rangle| \end{aligned}$$

$$\text{Prob}(2 | \psi_1) = 0$$

$$\begin{aligned} P(\text{conclusive}) &= P(1 | \psi_1) P(2 | \psi_1) + P(1 | \psi_2) P(1 | \psi_2) \\ &= \frac{1}{2} (2) (1 - |\langle \psi_1 | \psi_2 \rangle|) \end{aligned}$$

$$P_{\text{concl.}} = 1 - |\langle \psi_1 | \psi_2 \rangle|$$

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How similar are 2 pbits?

$$\cdot D(\vec{p}, \vec{q}) = \frac{1}{2} \sum |p_i - q_i|$$

$$\cdot F(\vec{p}, \vec{q}) = \sum \sqrt{p_i q_i}$$

How similar are 2 qubits?

$$\cdot D(\rho, \sigma) = \frac{1}{2} \text{Tr} |\rho - \sigma|$$

New today →

$$\cdot F(\rho, \sigma) = \text{tr} \sqrt{\rho^{1/2} \sigma \rho^{1/2}}$$

D features:

① $D(\rho, \sigma) = \frac{1}{2} \|\vec{r} - \vec{s}\|$ (qubits)

② D is a metric

③ $D(U\rho U^\dagger, U\sigma U^\dagger) = D(\rho, \sigma)$

④ $D = \max_{\{E_k\}} \text{tr}[E_k(\rho - \sigma)]$

⑤ $D(a(\rho), a(\sigma)) \leq D(\rho, \sigma)$ ("contractive")

⑥ $D(\sum p_i \rho_i, \sigma) \leq \sum p_i D(\rho_i, \sigma)$ ("convex")

Aside!

$D(\vec{p}, \vec{q})$ features

② D is a metric

③ $D(P_{\vec{p}}, P_{\vec{q}}) = D(\vec{p}, \vec{q})$

④ $D = \max_s \sum_{i \in S} |p_i - q_i|$

"max diff in probabilities as a sum will give"

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F features

- ① $F(p, \sigma) = F(\sigma, p)$
- ② $0 \leq F(p, \sigma) \leq 1$. $F(p, \sigma) = 1 \Leftrightarrow \sigma = p$
- ②b $F(p, \sigma) = 0 \Leftrightarrow p, \sigma$ have support on orthogonal spaces.
- ③ $F(U_p U^T, U_\sigma U^T) = F(p, \sigma)$
- ④ $F(p, \sigma) = \max_{\substack{|\psi\rangle: p = \frac{1}{2}(|\psi\rangle\langle\psi| + |\psi^\perp\rangle\langle\psi^\perp|) \\ |\phi\rangle: \sigma = \frac{1}{2}(|\phi\rangle\langle\phi| + |\phi^\perp\rangle\langle\phi^\perp|)}} |\langle\psi|\phi\rangle|$ [Uhlmann's Thm]
- ⑤ $F(p, \sigma) = \min_{\substack{P, Q \text{ perms} \\ \in \mathbb{R}^3}} F(\text{tr}(p E_k), \text{tr}(\sigma E_k))$
- ⑥ $F(p, |\psi\rangle\langle\psi|) = \sqrt{\langle\psi|p|\psi\rangle}$, $F(|\psi\rangle\langle\psi|, |\psi\rangle\langle\psi|) = |\langle\psi|\psi\rangle|$
- ⑦ $A(p, \sigma) = \arccos F(p, \sigma)$ is a metric
- ⑧ $F(a(p), a(\sigma)) \geq F(p, \sigma)$
- ⑨ $F(\sum p_i, \sigma) \geq \sum p_i F(p_i, \sigma)$ (F is concave)

Theorem:

$$1 - F(p, \sigma) \leq D(p, \sigma) \leq \sqrt{1 - F(p, \sigma)}$$

See NC § 9.2, 3 for proof

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Quantum Error Correcting Criteria:

Noise channel $\mathcal{N}: \rho \rightarrow \sum A_k \rho A_k^\dagger$, $\sum A_k^\dagger A_k = I$

e.g. Bit-flip channel on 3 qubits:

$A_0 = (1-p)^3 I \otimes I \otimes I$	$A_4 = (1-p)p^2 I X X$
$A_1 = (1-p)^2 p I I X$	$A_5 = (1-p)p^2 X I X$
$A_2 = (1-p)^2 p I X I$	$A_6 = (1-p)p^2 X X I$
$A_3 = (1-p)^2 p X I I$	$A_7 = p^3 X X X$

Correctible error channel $\mathcal{E}: \rho \rightarrow \sum E_k \rho E_k$, $\sum E_k^\dagger E_k \leq I$

e.g. $E_0 = A_0$, $E_2 = A_2$
 $E_1 = A_1$, $E_3 = A_3$

$[[n, k, 2t+1]]$ code
 when $\{E_i\}$ = all errors of
 "Pauli weight" $\leq t$

Desire recovery operator \mathcal{R} such that

e.g. Shor code is a $[[9, 1, 3]]$ code

\propto b/c \mathcal{E} is not TP

$\mathcal{R} \circ \mathcal{E}(\rho_{code}) \propto \rho_{code}$

e.g. ρ_{code} has basis $\{|i\rangle = |000\rangle, |j\rangle = |111\rangle\}$

Proof: Next time get this far

Thm: \mathcal{R} exists iff $\langle i | E_a^\dagger E_b | j \rangle = C_{ab} \delta_{ij}$, $C = C^\dagger$

Meaning:

- ① E_a on $|i\rangle$ distinguishable from E_b on $|j\rangle$: $\langle i | E_a^\dagger E_b | j \rangle = 0$ if $i \neq j$
- ② Measuring errors doesn't distinguish codewords: $\langle i | E_a^\dagger E_b | i \rangle = \langle j | E_a^\dagger E_b | j \rangle$