Phys 522 Quantum Mechanics II

Exam #3 (12.5% of course grade) Due Saturday, April 24 (at noon in TA's mailbox)

Spring 2010

This is an examination for Phys522. Your score will determine 12.5% of your grade for the course.

This is a take-home, open-book exam. You may use the textbook, your own notes, your own homework assignments, all class handouts, including solution sets for homework assignments, but do not consult books other than our primary textbook. You may use as much time as you want to complete the exam, as long as it is turned in by the deadline. Your completed exam must be solely your own work; do not consult anyone else in doing the exam.

Do your work on sheets of paper separate from the exam. The solution you hand in for a problem should be neat and legible and in a logical order. It should represent your best effort at solving the problem and should not include false starts and detours that led nowhere.

Staple your solution(s) to this page, and label this page with your name. Then *sign* and date the pledge below, and turn in the exam to the TA's mailbox. You may keep the exam question(s).

I have obeyed the rules in taking this exam. In particular, I have not consulted anyone about the exam, and the completed exam is solely my own work.

Signature

Date

Phys 522 Quantum Mechanics II

Exam #3

Spring 2010 (100 points)

Problem 1 (100 points) Consider a one-dimensional harmonic oscillator with mass m and frequency ω , i.e., a particle with Hamiltonian

$$H_0 = \frac{P^2}{2m} + \frac{1}{2}m\omega^2 X^2 = \hbar\omega(a^{\dagger}a + \frac{1}{2}) ,$$

where the annihilation operator is given by

$$a = \sqrt{\frac{m\omega}{2\hbar}} \left(X + i \frac{P}{m\omega} \right) \; .$$

The oscillator is initially, at time t = 0, in the ground state $|0\rangle$.

At time t = 0, a parametric interaction is turned on, which changes the Hamiltonian to

$$H = \frac{P^2}{2m} + \frac{1}{2}m\omega^2 [1 + g\cos(2\omega t)]X^2 .$$

The parametric interaction modulates the resonant frequency at twice the resonant frequency. The strength of the modulation is characterized by the dimensionless constant g. The parametric interaction can be treated as a perturbation,

$$W(t) = \frac{1}{2}gm\omega^2 X^2 \cos(2\omega t) ,$$

on top of the unperturbed Hamiltonian H_0 .

The state of the system at time t can be written in the energy-eigenstate (number) basis as

$$|\psi(t)\rangle = \sum_{n=0}^{\infty} c_n(t)|n\rangle = e^{-i\omega t/2} \sum_{n=0}^{\infty} b_n(t) e^{-in\omega t}|n\rangle .$$

In the following we are interested in the state of the oscillator after the parametric interaction has been on for a time τ . Throughout the problem we are only interested in the secular (resonant) terms. These secular terms do not oscillate at frequency ω , and they dominate when $\omega \tau \gg 1$.

(a) (25 points) Find the first-order secular contributions to the number-state amplitudes; i.e., find the secular terms in $b_n^{(1)}(\tau)$.

(b) (35 points) Find the second-order secular contributions to the number-state amplitudes; i.e., find the secular terms in $b_n^{(2)}(\tau)$.

(c) (40 points) Go to an interaction picture relative to H_0 . Find the interaction-picture Hamiltonian, make the rotating-wave approximation (i.e., keep only resonant terms), and integrate the result to find the interaction-picture evolution operator $U_I(\tau, 0)$. Use this

result to find the secular contributions to the amplitudes $b_n(t)$ to the same order as in parts (a) and (b). Your results should agree with those of parts (a) and (b).

$$[X,P] = i\hbar \qquad [a,a^{\dagger}] = 1$$

$$a|n\rangle = \sqrt{n}|n-1\rangle$$
 $a^{\dagger}|n\rangle = \sqrt{n+1}|n+1\rangle$ $a^{\dagger}a|n\rangle = n|n\rangle$

 $e^{i\theta a^{\dagger}a}ae^{-i\theta a^{\dagger}a} = ae^{-i\theta}$