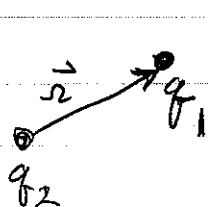


Physics 405 : Lecture 22

Introduction to Magnetostatics

Recall the definition of the electric field



$$\vec{F}_{12} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2} \hat{r} = \frac{1}{4\pi\epsilon_0} q_1 q_2 \left(\frac{\vec{r}}{r^3} \right)$$

Coulomb's law: Action at a distance

$$\Rightarrow \vec{F}_{12} = q_1 \vec{E}_2(\vec{r}_1)$$

local \vec{E} @ \vec{r}_1 due to q_2

$$\vec{E}_2(\vec{r}_1) = \frac{1}{4\pi\epsilon_0} q_2 \frac{(\vec{r}_1 - \vec{r}_2)}{|\vec{r}_1 - \vec{r}_2|^3}$$

Collection of
"external charges"
Static

Q "test charge" @ \vec{r}

$$\vec{F} = Q \vec{E}(\vec{r})$$

$$\vec{E} = \sum_i \frac{q_i}{4\pi\epsilon_0} \frac{\vec{r} - \vec{r}_i}{|\vec{r} - \vec{r}_i|^3}$$

$$= \int d^3r' \frac{\rho(\vec{r}')}{4\pi\epsilon_0} \frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3}$$

We took this empirical definition
and derived from it the general field eqns.
of electrostatics:

$$\boxed{\vec{\nabla} \cdot \vec{E} = \rho/\epsilon_0}, \quad \boxed{\vec{\nabla} \times \vec{E} = 0}$$

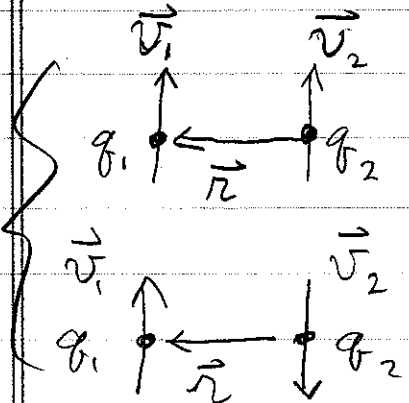
What happens when the charges move?

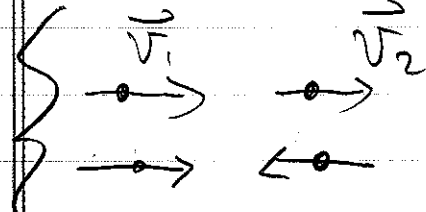
First empirical law: Constant velocity, $v \ll c$
(speed of light)

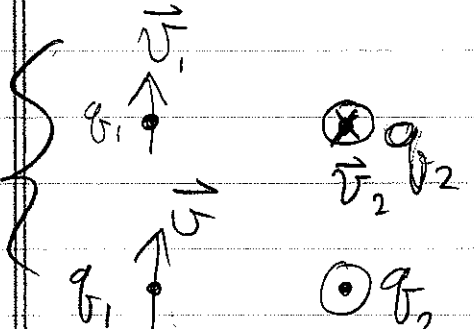
Consider two point charges, q_1, q_2 , separation vector \vec{r}
from $2 \rightarrow 1$, moving with velocities \vec{v}_1 and \vec{v}_2 .

• Empirical results: Additional force in addition to
coulomb

• Four cases

(i)  $\left. \begin{array}{l} \vec{v}_1 \uparrow \\ \vec{v}_2 \uparrow \end{array} \right\} \vec{F}_{12} \sim \mp \hat{r} \frac{q_1 q_2}{r^2} v_1 v_2$
on 1 due to 2

(ii)  $\left. \begin{array}{l} \vec{v}_1 \rightarrow \\ \vec{v}_2 \leftarrow \end{array} \right\} \vec{F}_{12} = 0$

(iii)  $\left. \begin{array}{l} \vec{v}_1 \uparrow \\ \vec{v}_2 \downarrow \end{array} \right\} \vec{F}_{12} = 0$

(iv)  $\left. \begin{array}{l} \vec{v}_1 \uparrow \\ \vec{v}_2 \leftarrow \end{array} \right\} \begin{array}{l} \vec{F}_{12} \sim \vec{v}_2 v_1 \frac{q_1 q_2}{r^2} \\ \vec{F}_{21} \sim 0 \end{array}$

The additional force between charges is a magnetic interaction

$$\vec{F}_{1,2}^{\text{magnetic}} = k_m \frac{q_1 q_2}{r^2} \vec{v}_1 \times (\vec{v}_2 \times \hat{r}) \quad (\text{non-relativistic})$$

Constant that determines the units

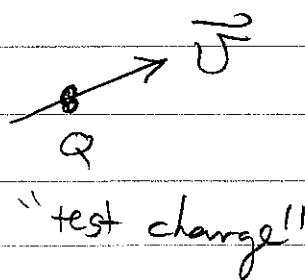
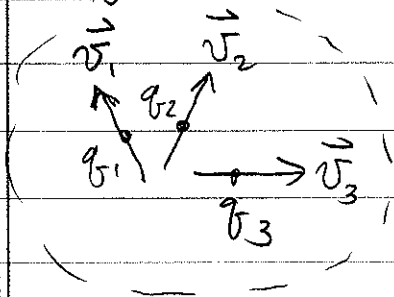
$$k_m = \frac{1}{4\pi\epsilon_0 c^2} = \frac{\mu_0}{4\pi}$$

$$\frac{\mu_0}{4\pi} = 10^{-7} \frac{(\text{Newtons})(\text{sec})^2}{(\text{Coulomb})^2}$$

Note: $|\vec{F}_{12}^{\text{mag}}| \ll |\vec{F}_{12}^{\text{electric}}|$

when $\frac{v}{c} \ll 1$ non-relativistic

Many moving charges



$$\vec{F}^{\text{magnetic}} = Q \vec{v} \times \vec{B}(\vec{r})$$

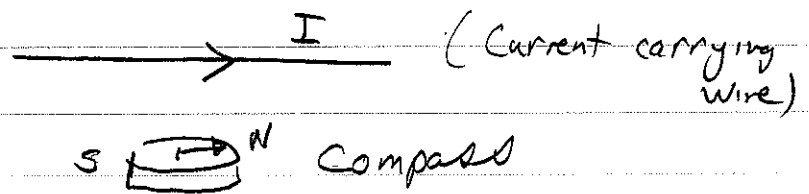
$$\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \sum_i q_i \vec{v}_i \times \frac{\hat{r}_i}{r_i^2}$$

Magnetic field

Conclusions

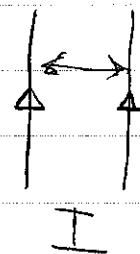
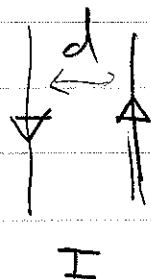
- (1) Charges in motion produce a magnetic field
Current is the source of \vec{B}

Oersted 1819

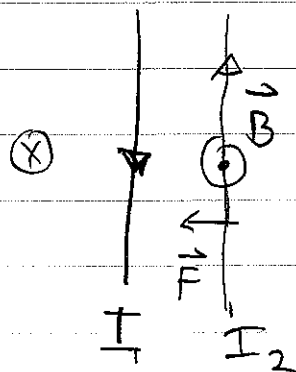


- (2) A Charge Q moving with velocity \vec{v} feels a force

$$\vec{F}_{\text{mag}} = Q \vec{v} \times \vec{B} \quad \text{Lorentz Force Law}$$



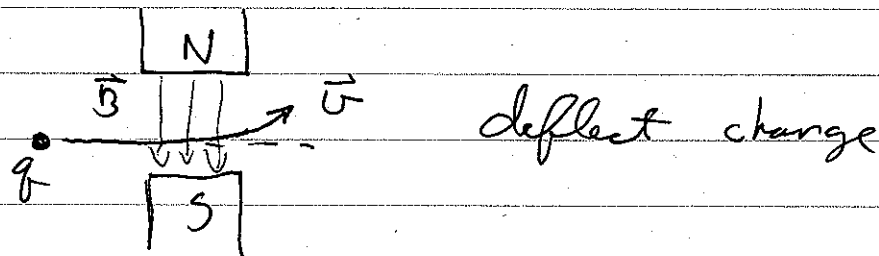
Anti-Parallel (parallel) current carrying
Wires attract (repell)



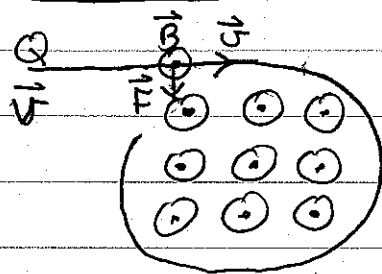
$$\text{Force/length} = \frac{\mu_0}{4\pi} \frac{I_1 I_2}{d}$$

Motion of charge particles in $\vec{E} + \vec{B}$ fields

Lorentz force $\vec{F} = Q(\vec{E} + \vec{v} \times \vec{B})$
 (usually \vec{E} dominates unless $|\vec{v}| \sim c$)



Cyclotron motion (Uniform \vec{B} field)



$$\vec{F}_m = Q\vec{v} \times \vec{B}$$

(Centripetal force)

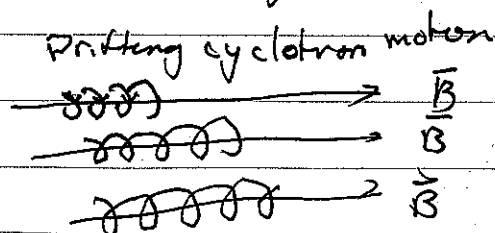
Circular motion $|\vec{F}_m| = QvB = Mv^2/R$

$$\Rightarrow R = \frac{Mv}{QB} \quad (\text{cyclotron radius})$$

identify momentum

Angular frequencies $\omega = \frac{v}{R} = \frac{QB}{M}$ (independent of v)

- Method of identify momenta of elementary particles
- Method of confining a plasma (gas of ions) for fusion



Currents: The Source of \vec{B} -field

$$I = \frac{\text{charge}}{\text{time}} \quad (\text{flow of charge})$$

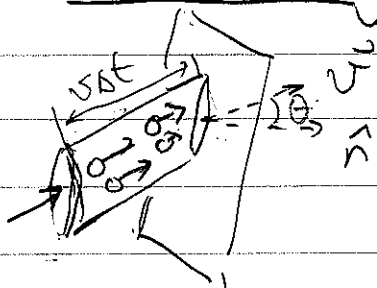
Consider fluid-like picture:

Charge density ρ , moving with local velocity \vec{v}

Define current density $\vec{J} = \rho(\vec{r}) \vec{v}(\vec{r})$

$$\Delta q = \rho(v \Delta t \Delta A \cos \theta)$$

$$\Rightarrow \frac{\Delta q}{\Delta t} = \int_S \vec{J} \cdot \hat{n} dA$$



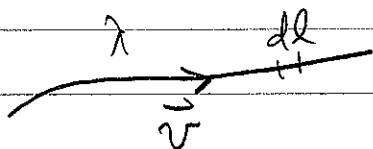
Flux of charge

$$I = \int_S \vec{J} \cdot d\vec{a}$$

\vec{J} is the source of \vec{B}
as ρ " " " of \vec{E}

Current Distributions (steady velocity)

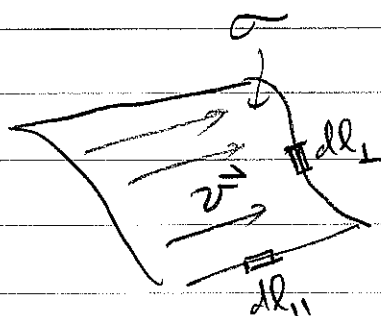
Line
charge
flow



$$dq = \lambda dl$$

$$\Rightarrow I = \lambda \frac{dl}{dt} = \lambda v$$

Surface
charge
flow



$$dq = \sigma da = \sigma dl_{\perp} dl_{\parallel}$$

$$\Rightarrow I = (\underbrace{\sigma v}_K) dl_{\perp}$$

$$[K] = \frac{\text{current}}{\text{length}}$$

$K \leftarrow$ surface current density

A hand-drawn diagram of a cylinder. The length of the cylinder is labeled as l with a double-headed arrow below it. The cross-sectional area of the cylinder is labeled as A with an arrow pointing to the right circular face.

$$\bar{J} = \rho v = \text{"Current density"}$$

$$[J] = \frac{\text{Current}}{\text{Area}}$$

Force on Continuous distribution

$$\vec{\Gamma} = \int dq \vec{v} \times \vec{B}$$

Line charge $dq = \lambda dl \Rightarrow F^L = \int I d\vec{l} \times \vec{B}$

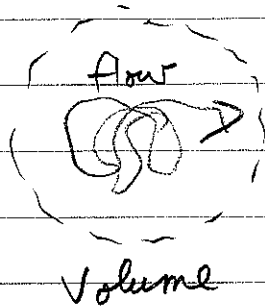
Surface charge $dq = \sigma da \Rightarrow \vec{F} = \int da \vec{K} \times \vec{B}$

Volume charge $dq = \rho d^3r \Rightarrow \vec{F} = \int d^3r \vec{J} \times \vec{B}$

Key is source of \vec{B}

$$\begin{array}{ccccccc} dq \vec{v} & \Leftrightarrow & I d\vec{\ell} & \Leftrightarrow & \vec{K} da & \Leftrightarrow & \vec{J} d^3x \\ \uparrow & & \uparrow & & \uparrow & & \uparrow \\ \text{point} & & \text{line} & & \text{surface} & & \text{volume} \end{array}$$

Conservation of Charge



Net Flow out of volume

$$I_{out} = \oint_S \vec{J} \cdot d\vec{a}$$

$$= - \frac{d}{dt} Q_{enc}$$

decrease in enclosed charge

$$\Rightarrow \oint_S \vec{J} \cdot d\vec{a} = - \frac{d}{dt} \int_V d^3r \rho(\vec{r}) = - \int_V d^3r \frac{\partial \rho}{\partial t}$$

$$\int (\vec{\nabla} \cdot \vec{J}) d^3r$$

$$\Rightarrow \boxed{\vec{\nabla} \cdot \vec{J} = - \frac{\partial \rho}{\partial t}}$$

'Magnetostatics' $\frac{\partial \rho}{\partial t} = 0$: No local change in charge density

$$\Rightarrow \boxed{\vec{\nabla} \cdot \vec{J} = 0} \quad \text{Magnetostatics}$$