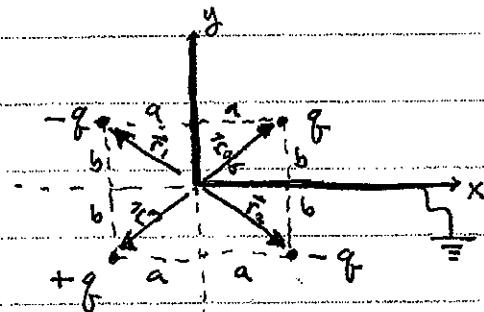


# Physics 405. P.S.#6 Solutions

## (1) Griffiths Problem 3.10



Finding the image charges is more of an art than a science. That is, one must "guess" the answer and then check it.

From the classic image problem of a point charge above a single plane we know that Laplace's Eq., and the boundary conditions are satisfied, so you might have chosen the 2 negative charges in the 2<sup>nd</sup> and 3<sup>rd</sup> quadrant.

However, you will find that the potential is not zero on the planes. We need a 3<sup>rd</sup> charge, +q, in the 3<sup>rd</sup> quadrant.

### Configuration of image charges

-q at  $\vec{r}_1 = -a\hat{x} + b\hat{y}$  and  $\vec{r}_2 = +a\hat{x} - b\hat{y}$

+q at  $\vec{r}_3 = -a\hat{x} - b\hat{y}$

The potential is that of the real charge plus 3 "images"

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \left( \frac{q}{|\vec{r}-\vec{r}_1|} - \frac{q}{|\vec{r}-\vec{r}_2|} - \frac{q}{|\vec{r}-\vec{r}_3|} + \frac{q}{|\vec{r}-\vec{r}_4|} \right)$$

$$\Rightarrow V(x, y, z) = \frac{q}{4\pi\epsilon_0} \left( \frac{1}{\sqrt{(x-a)^2 + (y-b)^2 + z^2}} - \frac{1}{\sqrt{(x+a)^2 + (y-b)^2 + z^2}} + \frac{1}{\sqrt{(x-a)^2 + (y+b)^2 + z^2}} + \frac{1}{\sqrt{(x+a)^2 + (y+b)^2 + z^2}} \right)$$

Check:

•  $\nabla^2 V = 0$  since we constructed it from such solutions

• Boundary conditions

(1) There is a charge at  $\vec{r}_q = a\hat{x} + b\hat{y}$  by construction

$$(2) V(x, 0, z) = \frac{q}{4\pi\epsilon_0} \left( \frac{1}{\sqrt{(x-a)^2 + y^2 + z^2}} - \frac{1}{\sqrt{(x+a)^2 + y^2 + z^2}} - \frac{1}{\sqrt{(x-a)^2 + y^2 + z^2}} + \frac{1}{\sqrt{(x+a)^2 + y^2 + z^2}} \right) = 0 \checkmark$$

(3)  $V(0, y, z) = 0 \checkmark$

The force on  $q$  is the electric field at  $q$  due to the surface charge on the conducting planes times  $q$ . But the  $E$ -field at  $\vec{r}_q$  due to the surface charge is exactly mimicked by the three image charges. Thus we can use Coulomb's law:  $\vec{F}_q = q \vec{E}(\vec{r}_q)$

$$\vec{F}(\vec{r}_q) = \frac{q}{4\pi\epsilon_0} \left\{ -\frac{q}{|\vec{r}_q - \vec{r}_1|^3} (\vec{r}_q - \vec{r}_1) + \frac{-q}{|\vec{r}_q - \vec{r}_2|^3} (\vec{r}_q - \vec{r}_2) + \frac{q}{|\vec{r}_q - \vec{r}_3|^3} (\vec{r}_q - \vec{r}_3) \right\}$$

Plugging in for  $\vec{r}_q, \vec{r}_1, \vec{r}_2, \vec{r}_3$  we get

$$\vec{F}(\vec{r}_q) = \frac{q^2}{16\pi\epsilon_0} \left\{ -\frac{\hat{x}}{a^2} - \frac{\hat{y}}{b^2} + \frac{a\hat{x} + b\hat{y}}{(a^2 + b^2)^{3/2}} \right\}$$

$$= -\frac{q^2}{16\pi\epsilon_0} \left\{ \left(1 - \frac{1}{(1 + \frac{b^2}{a^2})^{3/2}}\right) \frac{\hat{x}}{a^2} + \left(1 - \frac{1}{(1 + \frac{a^2}{b^2})^{3/2}}\right) \frac{\hat{y}}{b^2} \right\}$$

The work necessary to bring  $q$  in from  $\infty$  is just the ~~work backwards~~  $\rightarrow$  of the electrostatic energy of the 4 charge system since it takes no work to bring the image charges in from  $\infty$  (see handout given in class)

$$\begin{aligned} W &= \frac{1}{4} \left\{ \frac{1}{4\pi\epsilon_0} \frac{q q_1}{r_{q_1}} + \frac{1}{4\pi\epsilon_0} \frac{q q_2}{r_{q_2}} + \frac{1}{4\pi\epsilon_0} \frac{q q_3}{r_{q_3}} \right. \\ &\quad \left. + \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r_{q_2}} + \frac{1}{4\pi\epsilon_0} \frac{q_1 q_3}{r_{q_3}} + \frac{1}{4\pi\epsilon_0} \frac{q_2 q_3}{r_{q_3}} \right) \\ &= \frac{1}{4} \left( \sum_{\substack{i,j \\ i \neq j}} \frac{q_i q_j}{4\pi\epsilon_0 r_{q_i q_j}} \right) \end{aligned}$$

Plugging in for the charges and the distances between them

$$W = \frac{q^2}{16\pi\epsilon_0} \left\{ -\frac{1}{2a} - \frac{1}{2b} + \frac{1}{2\sqrt{a^2+b^2}} \right\} + \frac{1}{16\pi\epsilon_0} \left\{ \frac{-1}{a} - \frac{1}{b} + \frac{1}{\sqrt{a^2+b^2}} \right\}$$

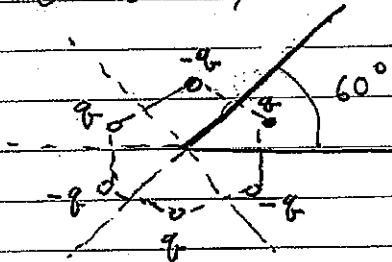
$\Rightarrow W = \frac{q^2}{16\pi\epsilon_0} \left\{ \frac{-1}{a} - \frac{1}{b} + \frac{1}{\sqrt{a^2+b^2}} \right\}$ 

 Check that  
 $W = -\vec{S}\vec{F} \cdot d\vec{l}$   
 where  $\vec{F}$  is the calculated force

The method will work if  $180^\circ$  is divisible by the angle between the planes.

$$\text{i.e. } \theta = \frac{180^\circ}{n} \quad n = 1, 2, 3, \dots$$

Example  $\theta = 60^\circ, n = 3$



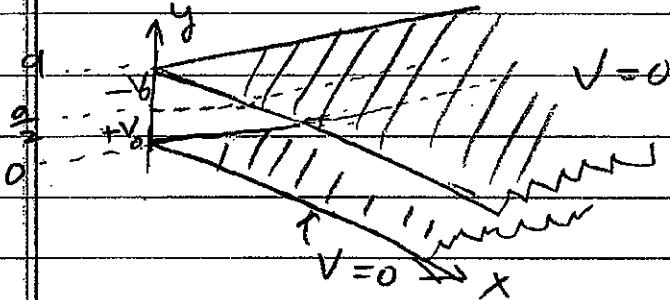
The solid circle is the true charge and the images are open circles

In general we need  $n-1$  image charges

As long as the charges sit on the corners of a polygon ~~such that~~ such that the conductivity planes bisect the sides, the boundary conditions are satisfied

Problem 3: Griffiths 3.12

Infinite slit with a step potential along  $y$



Same as example 3.1  
in Griffiths,

$$\text{Except: } V_0(y) = \begin{cases} V_0 & 0 \leq y < \frac{a}{2} \\ V_0 - \frac{aV_0}{2} & \frac{a}{2} \leq y < a \end{cases}$$

With the boundary conditions  $V=0$  at  $y=0$ ,  $y=a$ ,  $x=\infty$ ,  
the general solution of Laplace's eqn  $\nabla^2 V = 0$  is

$$V(x, y) = \sum_{n=1}^{\infty} C_n e^{-k_n x} \sin(k_n y)$$

$$k_n = \frac{n\pi}{a}$$

To find the expansion coefficients, we must match  $\partial_x V(x=0)$

$$V(0, y) = V_0(y) = \sum_{n=1}^{\infty} C_n \sin(k_n y)$$

"Project" onto the orthogonal function  $\sin(k_n y)$

$$\Rightarrow \int_0^a V_0(y) \sin(k_n y) dy = \sum_{n=1}^{\infty} C_n \underbrace{\int_0^a dy \sin(k_n y) \sin(k_n y)}_{\frac{a}{2} S_{nn}}$$

$$\Rightarrow C_n = \frac{2}{a} \int_0^a dy V_0(y) \sin(k_n y) =$$

$$= \frac{2V_0}{a} \left[ \int_0^{a/2} dy \sin\left(\frac{m\pi y}{a}\right) - \int_{a/2}^a dy \sin\left(\frac{m\pi y}{a}\right) \right]$$

(next Page)

$$\Rightarrow C_m = \frac{2}{\pi m} V_0 \left( -\cos\left(\frac{m\pi y}{a}\right) \Big|_0^{a_2} + \cos\left(\frac{m\pi y}{a}\right) \Big|_0^{a_2} \right)$$

$$= \frac{2}{\pi m} V_0 \left( 1 + \cos(m\pi) - \cos\left(\frac{m\pi}{2}\right) \right)$$

$$\Rightarrow C_m = \begin{cases} \frac{8V_0}{\pi m} & \text{if } m=2, 6, 10, 14 \text{ etc} \\ 0 & \text{otherwise} \end{cases}$$

(Mathematica, to follow)

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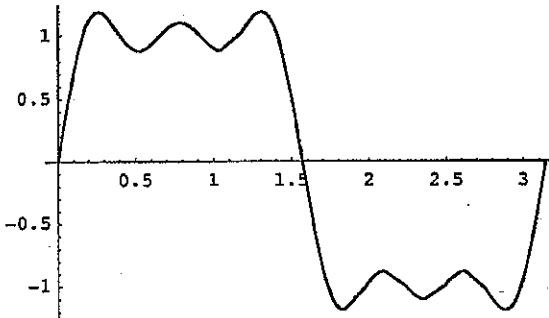
**Define the function**

```
phi[x_,y_,m_] := Exp[-m x] Sin[m y]
c[m_] := 2/(m Pi)*(1+Cos[m Pi] - 2 Cos[m Pi/2])
V[x_,y_,Ntot_] := Sum[c[m]*phi[x,y,m], {m,1,Ntot}]
```

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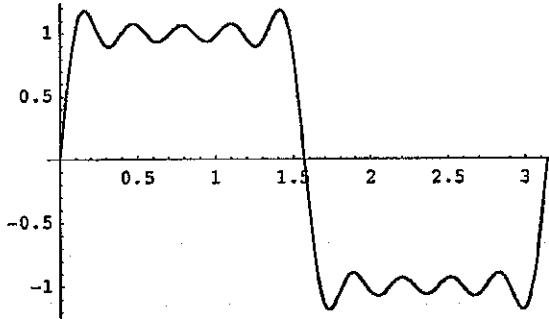
**Plots of the potential (in units of V0) along the y-axis for as Ntot gets larger.**

```
Plot[V[0,y,10],{y,0,Pi},PlotRange->All]
```



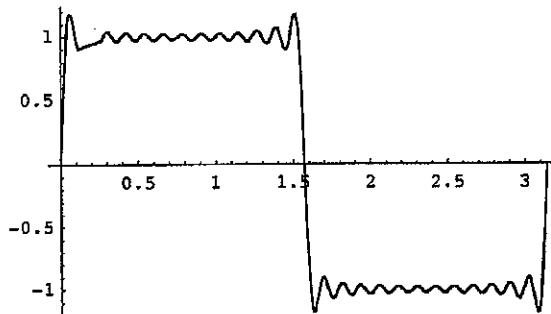
-Graphics-

```
Plot[V[0,y,20],{y,0,Pi},PlotRange->All]
```



-Graphics-

```
Plot[V[0,y,50],{y,0,Pi},PlotRange->All]
```

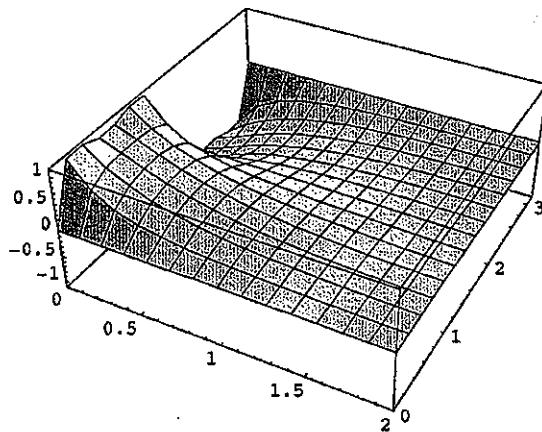


-Graphics-

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### 3D plot of the potential (in units of V0) for Ntot=20.

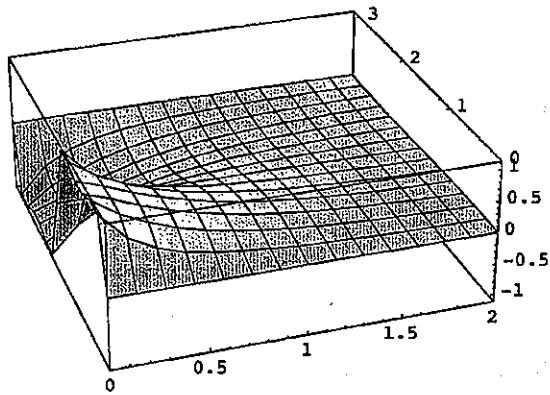
```
Plot3D[V[x,y,20],{x,0,2},{y,0,Pi},PlotRange->All]
```



-SurfaceGraphics-

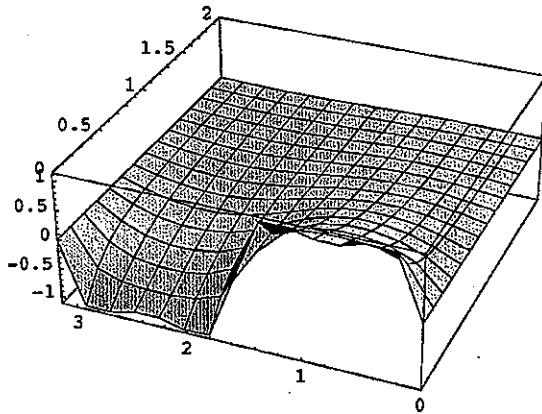
**■ different viewpoints**

```
Show[% , ViewPoint->{-1.533, -4.596, 2.517}]
```



-SurfaceGraphics-

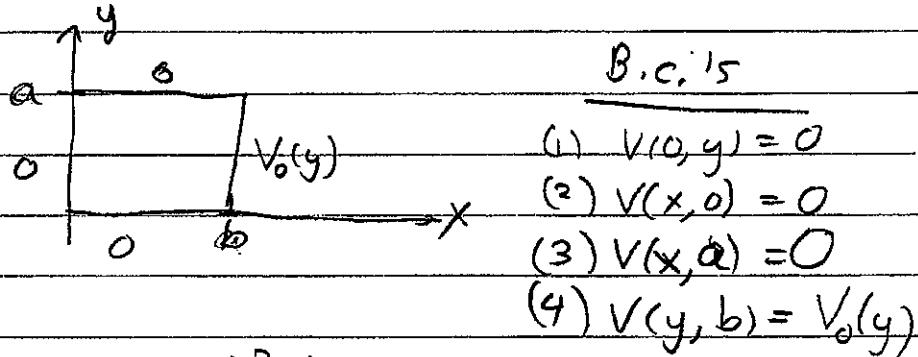
```
Show[% , ViewPoint->{-4.000, -1.670, 2.540}]
```



-SurfaceGraphics-

Problem 1: Griffiths 3.14

Rectangular pipe along  $z$ -axis with three sides grounded, and the fourth set  $\Rightarrow V_0(y)$



Separated solution to  $\nabla^2 V = 0$  in  $x$  and  $y$

$$(I) \quad V_k(x, y) = (A e^{kx} + B e^{-kx}) (C \sin ky + D \cos ky)$$

$$(II) \quad V_k(x, y) = (A e^{ky} + B e^{-ky}) (C \sin kx + D \cos kx)$$

In order to satisfy B.C.'s (2) and (3), the potential must go to zero in two position along  $y$

$\Rightarrow$  We must have an oscillating function of  $y$

$\Rightarrow$  The "normal modes" must be of type (I)

$$V_k(x, y) = (A e^{kx} + B e^{-kx}) (C \sin ky + D \cos ky)$$

$$\underline{\text{B.C. #1}} \Rightarrow V_k(0, y) = (A + B) (C \sin ky + D \cos ky) = 0$$

$$\Rightarrow A = -B$$

$$\Rightarrow V_k(x, y) = (C \sin ky + D \cos ky) \sinh(kx)$$

where  $\sinh(kx) = \frac{e^{kx} - e^{-kx}}{2}$ , and I absorbed  $2A$  into  $C + D$

$$\text{B.C. #2} \quad V_k(x, 0) = D \sin(kx) = 0$$

$$\Rightarrow D = 0$$

$$\text{B.C. #3} \quad V_k(x, a) = C_k \sinh(kx) \sin(ka) = 0$$

$$\Rightarrow \boxed{k = \frac{n\pi}{a}} \quad n = 1, 2, 3, \dots$$

$\Rightarrow$  General solution, satisfying b.c.'s # 1-3

$$\boxed{V(x, y) = \sum_{n=1}^{\infty} C_n \sinh(k_n x) \sin(k_n y)}$$

To find the coefficients "C<sub>n</sub>", we use the final b.c.

$$\text{B.C. #4} \quad V(b, y) = \sum_{n=1}^{\infty} C_n \sinh(k_n b) \sin(k_n y) = V_0(y)$$

$$\text{Use orthogonality of } \sin(k_n y): \int_0^a dy \sin(k_n y) \sin(k_m y) = \frac{a}{2} \delta_{mn}$$

$$\Rightarrow \boxed{C_n = \frac{2}{a \sinh(k_n b)} \int_0^a dy \sin(k_n y) V_0(y)}$$

(b) For  $V_0(y) = V_0$  constant

$$C_n = \frac{2V_0}{a \sinh(k_n b)} \left| \int_0^a dy \sin(k_n y) \right| \begin{cases} 0 & n \text{ even} \\ \frac{2a}{n\pi} & n \text{ odd} \end{cases}$$

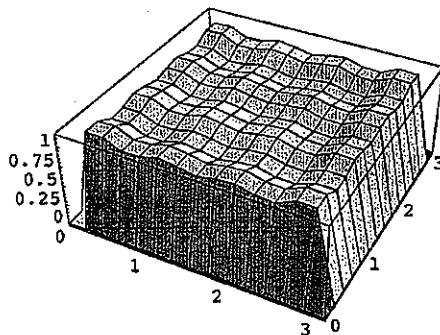
$$\Rightarrow \boxed{V(x, y) = \frac{4V_0}{\pi} \sum_{n=1}^{\infty} \frac{\sinh(\frac{n\pi}{a} x)}{n \sinh(\frac{n\pi}{a} b)} \sin(\frac{n\pi}{a} y)}$$

**Define the potential  $V(x,y,z)$**

```
v[x_,y_,z_,Ntot_] :=  
  (16/Pi^2)* Sum[1/(n m) * Sin[n x]*Sin[m y]*  
   (Sinh[(n^2+m^2)^.5 z]/Sinh[(n^2+m^2)^.5 Pi]),  
  {n,1,Ntot,2},{m,1,Ntot,2}]
```

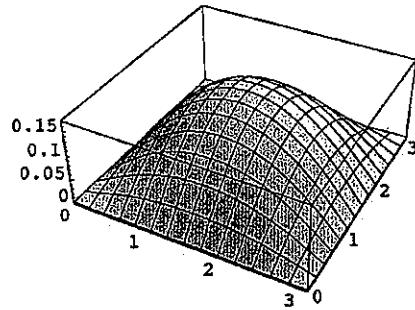
**Plots of the potential in the  $z=\pi$  plane  
(should approach zero).**

```
Plot3D[V[x,y,Pi,20],{x,0,Pi},{y,0,Pi},PlotRange->All]
```



**Plot of the potential in the  $z=\pi/2$  ( $x,y$ ) plane.**

```
Plot3D[V[x,y,Pi/2,10],{x,0,Pi},{y,0,Pi},PlotRange->All]
```

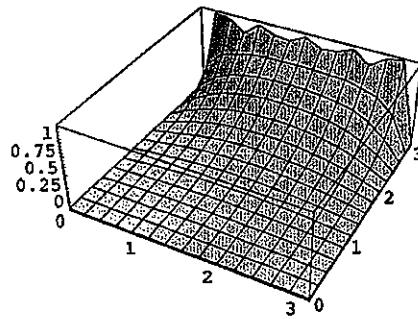


(\* Note that the boundary conditions are satisfied in x and y.  
This Plot seems to have a local maximum. How can this be? \*)

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**Plot of the potential in the  $y=\pi/2$  ( $x,z$ ) plane.**

```
Plot3D[V[x,Pi/2,z,10],{x,0,Pi},{z,0,Pi},PlotRange->All]
```

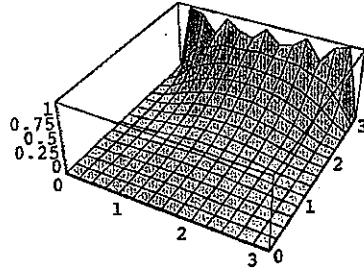


(\* Note that the boundary conditions are satisfied in  $x$  and  $z$ . \*)

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**Plot of the potential in the ( $x = y, z$ ) plane.**

```
Plot3D[V[x,x,z,10],{x,0,Pi},{z,0,Pi},PlotRange->All]
```



-SurfaceGraphics-

(\* This plot looks just like the previous one?  
Why? This about  $x,y$  symmetry \*)