

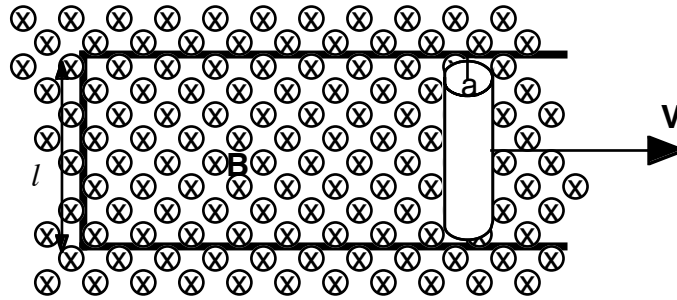
## Physics 511

### Problem Set #5: Faraday's Law, Quasistatics, and Conservation Laws DUE Fri. Mar. 10

Read Jackson 3rd Edition, 5.15-5.17, 6.1, 6.7  
Low, 2.5-2.6, 3.1-3.3

#### Problem 1: (10 Points)

A cylindrical resistive material with mass density  $\rho_m$  (mass/volume) and conductivity  $\sigma$  slides frictionlessly on two parallel conducting rails. The cylinder has length  $l$  and cross sectional area  $A$ . A uniform magnetic field  $\mathbf{B}$ , pointing into the page fills the entire region. The bar moves to the right, starting with a velocity  $v_0$ .



- (a) What current flows through the circuit, and in what direction?
- (b) What is the magnetic force on the rod? In what direction?
- (c) Use Newton's equation to show that velocity of the rod as function of time is

$$v(t) = v_0 e^{-\Gamma t}, \text{ where } \Gamma = \frac{B^2 \sigma}{\rho_M c^2}$$

- (d) The initial kinetic energy was  $\frac{1}{2} m v_0^2$ . At  $t = \infty$  the rod loses all this energy. Where does it go? Prove that energy is conserved by showing that the total energy that goes into this sink is  $\frac{1}{2} m v_0^2$ .

**Problem 2: Energy conservation and magnetic fields (15 points)**

To see explicitly why we put energy into creating a magnetic field, consider the work per unit length necessary to spin up a long cylindrical shell of charge so that it rotates about its axis at constant angular frequency  $\omega_0$ . Let the radius of the cylinder be  $R$ , the surface charge density be  $\sigma$ , so that  $\lambda=2\pi R\sigma$  is the total charge per unit length. After we spin the cylinder up, we have

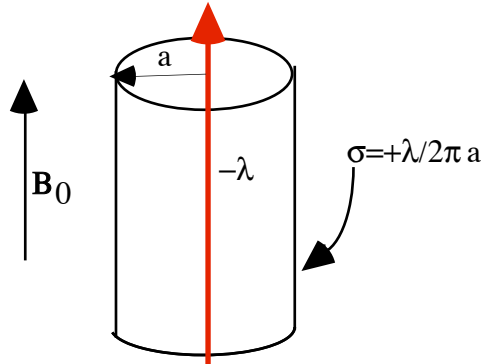
$$\mathbf{E}_0 = \begin{cases} 0, & \rho < R \\ \frac{2\lambda}{\rho} \hat{\mathbf{e}}_\rho, & \rho > R \end{cases}, \quad \mathbf{B}_0 = \begin{cases} \frac{4\pi}{c} \sigma \omega_0 R \hat{\mathbf{e}}_z, & \rho < R \\ 0, & \rho > R \end{cases}.$$

Now suppose we spin the cylinder up from rest, slowly over some time interval  $T$ . If  $T$  is much greater than  $R/c$ , we can neglect the displacement current. Then  $\mathbf{B}$  as a function of time is just the expression above with  $\omega_0$  replaced by  $\omega(t)$ . With  $\mathbf{B}(\mathbf{x},t)$  in hand, we can find the electric field as a function of time from Faraday's law:

- What is  $\mathbf{E}(\mathbf{x},t)$  as a function of time for all space? You may assume that the time dependent part of  $\mathbf{E}$  is in the  $\hat{\mathbf{e}}_\phi$  direction.
- If you are the external agent spinning up the shell, you must provide energy per unit volume at a rate  $-\mathbf{J} \cdot \mathbf{E}$ , over and above the energy you must put in to increase the mechanical rotational energy. Show that the volume integral of this expression per unit length along the  $z$ -axis is equal to the time rate of change of the total magnetic energy per unit length.
- From the above, it is clear that the "source" function for electromagnetic energy,  $-\mathbf{J} \cdot \mathbf{E}$ , is localized at  $\rho = R$ . Electromagnetic energy then flows from its creation point into  $\rho < R$ . Calculate the flow of energy per unit length to  $\rho < R$  by integrating the Poynting vector over a cylindrical surface just inside  $\rho = R$ . Show that the expression is equal to the time rate of change of the total magnetic energy per unit length inside  $\rho = R$ .
- You must also provide an additional torque per unit volume of  $-\mathbf{x} \times (\rho_{charge} \mathbf{E} + \frac{1}{c} \mathbf{J} \times \mathbf{B})$  to spin up the cylinder. Calculate the volume integral of this expression per unit length along the  $z$ -axis. Show that the rate of creation of electromagnetic angular momentum at  $\rho = R$  is just the energy creation rate divided by  $\omega$ . Comment on this result.
- Where does the electromagnetic angular momentum you are creating at  $\rho = R$  go to? In your answer consider the flux density of electromagnetic angular momentum,  $-\mathbf{x} \times (\vec{\mathbf{T}} \cdot \hat{\mathbf{n}})$ , just inside and outside of  $\rho = R$ . Where is this electromagnetic angular momentum stored? Can you get it back if you despin the cylinder?

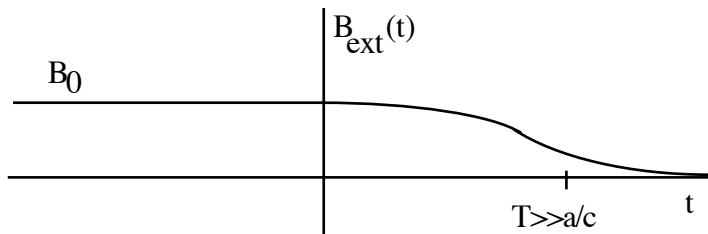
**Problem 3: Angular momentum conservation and magnetic fields (15 points)**

An infinitely long wire carries a charge per unit length of  $-\lambda$ , and lies along the  $z$ -axis. A plastic cylindrical shell of radius  $a$  is concentric about the  $z$ -axis and has a surface charge density  $\sigma = +\lambda / 2\pi a$  uniformly distributed over its surface (see sketch below).



The cylindrical shell is suspended in such a manner that it rotates freely about the  $z$ -axis. Its moment of inertia *per unit length* for rotation about its axis is  $I$ . The cylinder is initially at rest, and immersed in a uniform external magnetic field  $\mathbf{B}_0$ , produced by external currents.

- (a) Initially, what is the total electromagnetic angular momentum per unit length?
- (b) Now assume at  $t=0$  we slowly begin to reduce the external currents and  $B$ -field to zero over some time  $T \gg a / c$ ,  $\mathbf{B}_{\text{external}}(t) = B(t) \hat{\mathbf{e}}_z$



The cylinder will begin to rotate about its axis with angular frequency  $\omega(t)$ . What is the total magnetic field for  $\rho < a$  ?

- (c) From mechanics we know that the time rate of change of angular momentum equals the applied torque (per unit length)  $\frac{d}{dt}(I\omega) = (\mathbf{x} \times \lambda \mathbf{E})_z$ . Use this relation and Faraday's law to find an expression for  $d\omega / dt$
- (d) What is the final value of  $\omega$  (for  $t \gg T$ )? What is the final value of the magnetic field for  $\rho < a$ . Show that the final angular momentum of the system (mechanical plus electromagnetic) for  $t > T$  is the same as your part (a) for  $t < 0$ .