

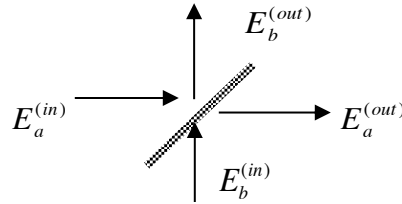
Physics 566, Quantum Optics

Problem Set #8

Due: Tuesday Oct. 28, 2025

Problem1: The beam splitter and other linear transformations (25 points)

Consider a symmetric beam splitter



In the first weeks of lecture, we showed that the pair $(E_a^{(out)}, E_b^{(out)})$ is related to $(E_a^{(in)}, E_b^{(in)})$ through a unitary “scattering matrix”

$$\begin{bmatrix} E_a^{(out)} \\ E_b^{(out)} \end{bmatrix} = \begin{bmatrix} t & r \\ r & t \end{bmatrix} \begin{bmatrix} E_a^{(in)} \\ E_b^{(in)} \end{bmatrix}$$

where $|t|^2 + |r|^2 = 1$, $\text{Arg}(t) = \text{Arg}(r) \pm \frac{\pi}{2}$, so that a possible transformation is,

$$E_a^{(out)} = \sqrt{T} E_a^{(in)} + i\sqrt{1-T} E_b^{(in)}, \quad E_b^{(out)} = \sqrt{T} E_b^{(in)} + i\sqrt{1-T} E_a^{(in)}, \quad \text{where } T = |t|^2.$$

Classically, if we inject a field only into one input port, leaving the other empty, the field in that mode will become attenuated, e.g., $E_a^{(out)} = \sqrt{T} E_a^{(in)} < E_a^{(in)}$.

(a) Consider now the quantized theory for these two modes, $E_a \Rightarrow \hat{a}$, $E_b \Rightarrow \hat{b}$. Suppose again that a field is injected only into the “a-port”. Show that

$$\hat{a}^{(out)} = \sqrt{T} \hat{a}^{(in)} \text{ is inconsistent with the quantum uncertainty.}$$

(b) In order to preserve the proper commutation relations we cannot ignore *vacuum fluctuations* entering the unused port. Show that if the “in” and “out” creation operators are related by the scattering matrix,

$$\begin{bmatrix} \hat{a}^{(out)\dagger} \\ \hat{b}^{(out)\dagger} \end{bmatrix} = \begin{bmatrix} t & r \\ r & t \end{bmatrix} \begin{bmatrix} \hat{a}^{(in)\dagger} \\ \hat{b}^{(in)\dagger} \end{bmatrix}, \text{ the commutator is preserved.}$$

(c) Suppose a single photon is injected into the a-port, so that the “in-state” is $|\psi^{(in)}\rangle = |1\rangle_a \otimes |0\rangle_b$. The “out-state” is $|\psi^{(out)}\rangle = \hat{S}|\psi^{(in)}\rangle$ where $\hat{S}$ is the “scattering operator”, defined so that $\hat{S}\hat{a}^{(in)\dagger}\hat{S}^\dagger = \hat{a}^{(out)\dagger}$ and $\hat{S}\hat{b}^{(in)\dagger}\hat{S}^\dagger = \hat{b}^{(out)\dagger}$.

$$\text{Show that } |\psi^{(out)}\rangle = t|1\rangle_a \otimes |0\rangle_b + r|0\rangle_a \otimes |1\rangle_b.$$

(d) Suppose a coherent state is injected into the a-port $|\psi^{(in)}\rangle = |\alpha\rangle_a \otimes |0\rangle_b$. Which is the output, $|\psi^{(out)}\rangle = |t\alpha\rangle_a \otimes |r\alpha\rangle_b$ or $|\psi^{(out)}\rangle = t|\alpha\rangle_a \otimes |0\rangle_b + r|0\rangle_a \otimes |\alpha\rangle_b$? Explain the difference between these.

(e) A general linear optical system consisting, e.g., of beam-splitters, phase shifters, mirrors, etalons, etc. can be described by a unitary transformation on the modes

$$E_k^{(out)} = \sum_{k'} u_{kk'} E_{k'}^{(in)}.$$

In the quantum description the mode operators transform by the scattering transformation

$$\hat{a}_k^{(out)\dagger} = \hat{S}\hat{a}_k^{(in)}\hat{S}^\dagger = \sum_{k'} u_{kk'}^T \hat{a}_{k'}^{(in)\dagger}, \text{ where } u_{kk'} \text{ is a unitary matrix.}$$

Show that if we start with a multimode coherent state $|\psi^{(in)}\rangle = |\{\alpha_k^{(in)}\}\rangle$, the output state is also a coherent state, $|\psi^{(out)}\rangle = |\{\alpha_k^{(out)}\}\rangle$, with $\alpha_k^{(out)} = \sum_{k'} u_{kk'} \alpha_{k'}^{(in)}$.

(f) The previous part highlights how linear transformations are essentially classical. This was true for input with exactly one photon or for coherent states. However, this is not true for more general inputs. Suppose we send one photon into *both ports*, of a 50-50 beam-splitter $T=1/2$, $|\psi^{(in)}\rangle = |1\rangle_a \otimes |1\rangle_b$. Show that the output state is,

$$|\psi^{(out)}\rangle = \frac{1}{\sqrt{2}}(|2\rangle_a |0\rangle_b + |0\rangle_a |2\rangle_b).$$

This says that the two photons both going to port-a or to port-b, but never one in port-a and one in port-b. This is an effect of Bose-Einstein quantum statistics. Explain in terms of destructive interference between indistinguishable processes.

Problem 2: Boson Algebra (25 points)

This problem is to give you some practice manipulating the boson algebra.

(a) Prove the (over) completeness integral for coherent states

$$\int \frac{d^2\alpha}{\pi} |\alpha\rangle\langle\alpha| = \hat{1} \quad (\text{Hint: Expand in number states}).$$

This basis is over-complete since as the coherent states are not orthonormal (see next part).

(b) Prove the group property of the displacement operator

$$\hat{D}(\alpha)\hat{D}(\beta) = \hat{D}(\alpha + \beta)\exp\{i\text{Im}(\alpha\beta^*)\}$$

$$\text{and thus } \langle\alpha|\beta\rangle = e^{-\frac{|\alpha-\beta|^2}{2}} e^{-i\text{Im}(\alpha\beta^*)}$$

(c) Show that the displacement operator has the following matrix elements

$$\text{Vacuum: } \langle 0|\hat{D}(\alpha)|0\rangle = e^{-|\alpha|^2/2}$$

$$\text{Coherent states: } \langle\alpha_1|\hat{D}(\alpha)|\alpha_2\rangle = e^{-|\alpha+\alpha_2-\alpha_1|^2/2} e^{i\text{Im}(\alpha\alpha_2^*-\alpha_1\alpha_2^*-\alpha_1\alpha_2^*)}$$

$$\text{Fock states: } \langle n|\hat{D}(\alpha)|n\rangle = e^{-|\alpha|^2/2} \mathbf{L}_n(|\alpha|^2), \text{ where } \mathbf{L}_n \text{ is the Laguerre polynomial of order } n$$

$$L_n(x) = \sum_{m=0}^n \binom{n}{m} \frac{(-1)^m}{m!} x^m$$

(d) Show that for a coherent state, the probability amplitude for detecting a quadrature is

$$\Psi(X) = \langle X|\alpha_0\rangle = e^{iP_0X} \frac{1}{\pi^{1/4}} e^{-\frac{(X-X_0)^2}{2}}, \quad \tilde{\Psi}(P) = \langle P|\alpha_0\rangle = e^{-iPX_0} \frac{1}{\pi^{1/4}} e^{-\frac{(P-P_0)^2}{2}},$$

where $\alpha_0 = \frac{X_0 + iP_0}{\sqrt{2}}$. Interpret this based on your understanding of wave mechanics.